

Cell Bulging and Extrusion in a Three-Dimensional Bubbly Vertex Model for Curved Epithelial Sheets

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Cell extrusion is an essential mechanism for controlling cell density in epithelial tissues. Another essential element of epithelia is curvature, which is required to achieve complex shapes, like in the lung or intestine. Here, we introduce a three-dimensional bubbly vertex model to study the interplay between extrusion and curvature. We find a generic cellular bulging instability at topological defects, which is much stronger than for standard vertex models. Analyzing cell shapes in three-dimensional imaging data of spherical mouse colon organoids, we infer that pentagonal cells have an increased basal interfacial tension, suggesting that cells at topological defects react to the different force conditions. Using the bubbly vertex model, we show that such basal tensions stabilize against the predicted instability and result in better cell shape control than tissue-scale mechanisms such as lumen pressure and spontaneous curvature. Our theory suggests that epithelial curvature naturally leads to bulged and extrusionlike cell shapes because the interfacial curvature of individual cells at the defects strongly amplifies buckling effected by tissue-scale topological defects in elastic sheets. Our results highlight the complex interplay of forces across scales in three-dimensional tissue organization.

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I. INTRODUCTION

Epithelial tissues are close-packed and sealed cell layers that form the barriers between different compartments in the body, for example, in the lung or the intestine. In order to regulate their cell density, they have to achieve a balance between cell division and cell death. Recently, it has become clear that both processes are strongly connected to mechanics: cells divide more if they are stretched by neighboring cells at low cell density, and they stop dividing and die if the epithelium becomes too crowded [1–3]. In order to make sure that cells leave the tissue without disrupting its barrier function, they are usually expelled by extrusion, which means they are squeezed out of the

epithelium without the generation of any voids [4,5]. Extrusion is especially prominent in the intestine, where the epithelium renews itself every 3–5 days [6]. Here, cells proliferate in the crypt regions and then move in a constant stream to the tip of villi, where they are extruded [7]. Very recently, it has been shown that most cells extruded in the villus regions are still alive, which means that mechanically they do not differ much from their neighbors [8]. The exact mechanisms of extrusion are not clear, but one important element seems to be the presence of specialized actin structures on the lateral and basal sides [3,8]. For flat model epithelia, it has been shown that cell extrusion occurs preferentially at topological defects in the nematic orientation field of the cells [9–11].

In general, cells in nonstratified epithelia tend to have hexagonal order because this is the close-packed optimum in two dimensions, similar to the situation in sphere packings and foams. Hexagonal order is also dominant in curved epithelial sheets [12–15]. However, now some pentagonal defects are required to generate curvature [13]. Because of Euler's polyhedron theorem, 12 pentagonal

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defects are required to achieve the topology of a sphere, with the possibility of additional pentagonal-heptagonal pairs. Thus, topological defects occur naturally in curved epithelia, while in flat cell layers, they arise rather in a stochastic manner from the movement of cells. Because epithelia are often curved due to their biological function, for example, in the crypts and villi of the intestine, the question arises as to whether they might use the naturally occurring topological defects to promote cell extrusion. While previous theoretical models have addressed such extrusion in a flat configuration focusing on fluidlike epithelia [16] or neglecting the three-dimensional nature of the problem [17], the coupling of curvature, defects, and shape instabilities has hardly been addressed before. Defects in such a setting have been linked to extrusion in flat epithelia on a substrate in phase field models of tissues [16] and to shape instabilities in vertex models [18], but the underlying mechanisms and their relation to curvature have not been addressed yet.

Recently, we have reported on a morphological instability at topological disclination defects in curved hexagonal epithelia in the three-dimensional vertex model (VM) [19]. This buckling instability arises because spherical epithelia are effectively elastic shells, similar to viral capsids [20,21] or buckyballs [22], for which such an instability has been known since the pioneering work by Seung and Nelson [23]. However, the standard VM does not consider that cells might adapt their own interfacial curvature to their mechanical environment. Here, we investigate this aspect by using the theoretical bubbly vertex model (BVM) and by analyzing experimentally observed cell shapes in spherical mouse colon organoids. In contrast to the normal VM, which describes multicellular systems with flat interfaces, in the BVM the interfaces are curved [24–28]. Such interfacial curvature leverages the three-dimensional nature of the cellular system and allows for different mechanisms than the usual VM approach to epithelia. On the experimental side, colon organoids are particularly suited to address the coupling of mechanics, cell shape, and topology, as they form spherical epithelial monolayers. They are derived from stem cells gathered from colonic crypts; they recapitulate developmental processes of the colon and thus serve as an important model system for development and disease [29].

We start in Sec. II by showing, with computer simulations, that the curved interfaces in the BVM strongly amplify the buckling at defects in curved sheets if compared to the standard VM with flat interfaces. Through segmentation of 3D microscopy data of mouse colon organoids with the subsequent inference of interfacial tensions, we find that cellular shapes are more regular in spherical organoids than theoretically expected. We also infer that outer (basal) tension is increased at pentagonal cells, suggesting that cells react to the changed force conditions at topological defects. In the BVM, we find

that such increased basal tensions stabilize against the predicted instability and yield full control over cellular bulging. To understand the underlying instability theoretically, we consider the representative case of a shell with icosahedral symmetry in which interfacial curvature leads to a significant lowering of a previously described icosahedral buckling threshold [19,20,23]. In Sec. III, we develop a continuum description of a defect cell in a mean-field tissue, allowing us to extend our model to describe and study extrusionlike cell geometries. We find that the possibility of curved interfaces results in a local decrease of the saddle splay modulus, thus leading to localization of Gaussian curvature in the defect. We find that lowering the apicobasal surface tension and a supra-cellular actin ring may initiate bulging, while only the former allows for a transition of a partially extruded geometry toward full extrusion. Our model predicts a lower energy barrier for extrusion at topological (pentagonal) defects. We conclude in Sec. IV with a discussion of our results.

II. BUBBLY VERTEX MODEL AND BULGING INSTABILITY AT TOPOLOGICAL DEFECTS

A. Definition of bubbly vertex model

For the description of tissue mechanics and shape, we adopt the usual tension-based approach. We start from the regular three-dimensional vertex model for epithelial monolayers in which individual cells are described as bounded volumes with one apical, one basal, and multiple lateral faces as boundaries. While in the classical vertex model these boundary faces are assumed to be flat by triangulating them with a passive interpolation point [19,30], we explicitly consider the effect of possible interface curvature on the mechanics and thus introduce the bubbly vertex model, in which the interfaces are allowed to curve [24–28], cf. Fig. 1(a).

Each cell i has one apical area $A_{a,i}$, one basal area $A_{b,i}$, and an integrated lateral area $A_{l,i}$, corresponding to the areas of the boundary faces. We assume the faces to have surface tensions Γ_a , Γ_b , and Γ_l for the apical, basal, and lateral faces, respectively, and that their values are the same for all cells (no dependence on i). The total energy of the tissue is then described by the sum of the individual surface area contributions,

$$E = \sum_i \left(\Gamma_a A_{a,i} + \Gamma_b A_{b,i} + \frac{1}{2} \Gamma_l A_{l,i} \right), \quad (1)$$

where the factor $1/2$ accounts for double counting. The apical and basal network topologies are identical; i.e., both apical and basal sides have the same number of adjacent apical and basal faces, respectively, and lateral faces are quadrilateral. We assume each cell has a fixed volume V_{cell} . We nondimensionalize the surface tension

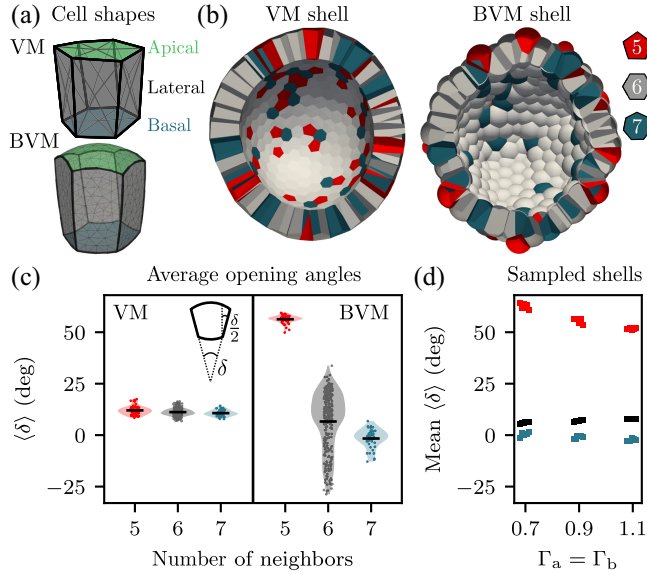


FIG. 1. Cell bulging at topological defects in the bubbly vertex model. (a) Schematic depiction of the VM representation of a cell with apical, basal, and lateral faces via an interpolation point of the polygonal faces, and of the BVM representation with curved interfaces. (b) Spherical VM and BVM shells after energy minimization with random initialization of cell centers. The BVM shell is after minimization in the Surface Evolver. The number of cell neighbors is color-coded. (c) Opening angles of the cells by neighbor number for the VM and BVM, where the half angle is determined for each lateral interface (inset) and then averaged over lateral interfaces. Violin plots show the distribution, with the horizontal line marking the mean value. The tensions are $\Gamma_a = \Gamma_b = 0.9$ and $\Gamma_1 = 1$. (d) Mean opening angles of cells with different neighbor numbers (color-coded; red pentagons, black hexagons, blue heptagons) in different BVM shells with random initial cell configurations ($N = 6$) and varying apical and basal tensions $\Gamma_a = \Gamma_b$ and $\Gamma_1 = 1$. All shells have 400 cells.

via Γ_1 and the volume via V_{cell} ; i.e., we set $\Gamma_1 = 1$ and $V_{\text{cell}} = 1$.

Because of the fluid viscoelastic behavior of the cell cortex [31,32], we neglect elastic in-plane contributions of the interfaces, as these will relax over sufficiently long time-scales. For now, we also do not consider a bending energy contribution [26], making use of the analogy to soap films, which are purely driven by surface tensions. This approach can be justified by the small thickness of the cortex, which will lead to small bending rigidities in the case of large interfacial tensions. We will relax this assumption later.

To differentiate between the model with and without membrane curvature, we call the former the VM and the latter the BVM throughout the rest of this article.

B. Instability at topological defects in the BVM

We consider a spherical epithelium with a randomized starting configuration. For simulations we use *OrganoidChaste* [33], a custom-written package for

Chaste [34] implementing a three-dimensional vertex model for monolayers, which was used in Ref. [19] and is described in Appendix A. Cell centers are created using the random sequential adsorption mechanism [35] and then a Voronoi tessellation on the sphere to determine cell shape. Using a rate of active topological T1 transitions, which is slowly reduced to zero, the energy is minimized to obtain a spherical starting configuration. This process is performed in the VM to obtain shapes with low energy, and a representative minimal energy configuration is shown in Fig. 1(b). This topology is then fixed for the comparison between VM and BVM shapes, as here we do not consider tissue dynamics on a larger timescale.

Next, we minimize the energy in the BVM by using the VM configuration as the starting configuration in the standard software Surface Evolver for energy minimization of surfaces under tension [36]. Again, we minimize at constant cell volume and obtain the result shown in Fig. 1(b). As shown by the color-coding, for both the VM and the BVM, pentagons, hexagons, and heptagons are distributed all over the surface. The spherical configuration necessitates curvature screening [37] through additional pairs of topological defects (pentagon-heptagon pairs) compared to an icosahedral configuration of a hexagonal lattice with 12 pentagons. While topological disclination defects indeed seem to be screened by the distribution of additional defect pairs in the VM, leading to an approximately spherical configuration, we observe an interesting instability in the BVM. Pentagonal cells tend to bulge outward, minimizing their luminal interfacial area, despite the fact that we do not observe a tissue-scale buckling instability. Figure 1(c) demonstrates the dramatic difference between the two models: The opening angles of pentagonal cells are much larger in the BVM than in the VM, agreeing with the visual impression when comparing shells in Fig. 1(b). In addition, we also see that the distribution of opening angles in the BVM is much broader, with some cells even bending in the opposite direction. We conclude that the BVM allows for more adaptation and variability in cell shape than the classical VM, with a particularly strong effect on the pentagonal defects. We verify the generality of this result by sampling many random configurations for shells with the same number of cells and for different apicobasal tensions $\Gamma_a = \Gamma_b$, effectively considering different shell sizes with respect to the monolayer thickness, cf. Fig. 1(d).

C. Cell shapes at topological defects in mouse colon organoids

Mouse colon (mCol) organoids are spherical epithelia derived from single stem cells [29,38]. Figure 2(a) shows a typical microscopy image. We have grown these organoids in extracellular matrix (basement membrane extract) from mechanically dissociated (split) closed organoids and fixed, stained, and imaged them at day 5 postsplitting. The luminal (inner) side of the organoids corresponds to the

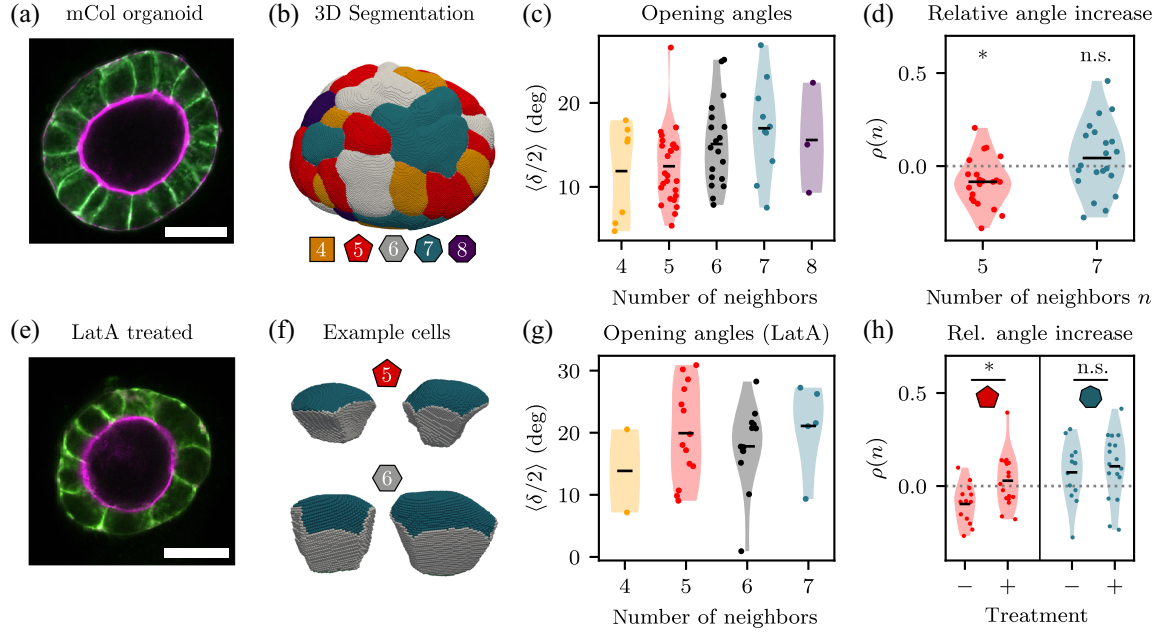


FIG. 2. Experimentally observed cell shapes at topological defects in mouse colon (mCol) organoids. (a) Microscopy image of spherical mouse colon organoid. Green is the membrane (β -catenin), and magenta is the F-actin. The luminal side is apical (high actin), and the outer side is basal (facing the matrix). (b) Segmented organoid in 3D and the number of neighbors (color-coded) determined from the basal polygonal edges. (c) Opening angles $\delta/2$, averaged over all lateral interfaces, which vary for different neighbor numbers. Each dot represents one cell, and the bars depict the means. (d) Relative angle increase at pentagonal and heptagonal cells, compared to hexagonal cells. Every data point is one organoid ($N = 22$). (e) Microscopy image of organoid treated with Latrunculin A (LatA). (f) Examples of pentagonal and hexagonal cells as segmented in 3D. (g) Opening angles for organoid treated with LatA. (h) Comparison of the relative angle increase ρ from wild-type ($N = 13$) and LatA-treated samples ($N = 17$), with larger effects at topological defects with LatA. Scale bars in panels (a) and (e) are $20\ \mu\text{m}$. Stars indicate the following: $*p < 0.05$, $**p < 0.01$, $***p < 0.001$. Tests against 0 in panel (d) are with Wilcoxon, and tests against each other in panel (h) are with Mann-Whitney-U tests.

apical side of the cells, as demonstrated by the actin staining in Fig. 2(a). The outer side corresponds to the basal side of the cells and has been growing against the matrix. The three-dimensional images were segmented, as shown in Fig. 2(b), using a previously described pipeline that we have developed in the context of force inference [39] (see Appendixes B and C for details). Similarly to the BVM, we determined the (half) cellular opening angles $\delta/2$ of the mCol cells as a function of their number of neighbors. Contrary to what we expect from the BVM, we find that, in some organoids, the opening angles of pentagonal cells are smaller than for hexagons and vice versa for heptagons [Fig. 2(c) for one organoid]. To allow for a comparison across organoids, we determined the relative increase of cell-averaged pentagonal ($n = 5$) and heptagonal ($n = 7$) opening angles, $\langle \delta/2 \rangle_n$, compared to the angles averaged over hexagonal cells, $\langle \delta/2 \rangle_6$, i.e.

$$\rho(n) = \frac{\langle \delta/2 \rangle_n - \langle \delta/2 \rangle_6}{\langle \delta/2 \rangle_6}. \quad (2)$$

The result is shown in Fig. 2(d). While we find this angle increase in pentagons to be larger than zero in some organoids, consistent with the bulging instability, it is

significantly lower than zero on the level of many organoids. In heptagonal cells, this difference is not significant.

The experimental data show that cell shape is indeed significantly different at topological defects, albeit oppositely to the bulging instability predicted by the BVM. As the actomyosin system has a stabilizing function through both the formation of an apical (luminal) actin belt and interfacial contractility, we treated the mCol organoids with Latrunculin A, an inhibitor of actin polymerization, effectively leading to actin depolymerization. While overall organoid morphology was not disrupted [Fig. 2(e)], we found a signature of the bulging instability in multiple treated organoids [Fig. 2(f) and 2(g)]. Indeed, looking at the relative angle increase $\rho(n)$, we found increased angles at pentagonal cells in treated organoid populations, hinting at a potentially stabilizing role of the actomyosin system against the bulging instability at disclination defects [Fig. 2(h)]. In heptagonal cells, we do not see a significant difference between treated and untreated organoids, which could be explained by the smaller relevance of this instability here due to the lack of saddlelike structures in these spherical organoids, cf. Fig. 1(c).

Cellular shapes display more regular opening angles than predicted by the BVM, suggesting more controlled force generation. In particular, mechanosensing and an active

compensation mechanism may stabilize the shape. In order to test this idea, we determined the interfacial tensions of the individual apical, basal, and lateral interfaces, denoted by $\hat{\Gamma}$, from the microscopy images. Assuming mechanical equilibrium, interfacial tensions lead to forces normal to tri-interfacial junctions [39–43]. Using our recently developed force inference method [39], we inferred the underlying relative tensions from the interfacial dihedral angles in the segmented images [cf. Fig. 3(a), see Appendix B for details]. Figures 3(b) and 3(c) display the different average tensions of different interface types at pentagonal and heptagonal cells, $\langle \hat{\Gamma} \rangle_n$, from the respective averages in hexagonal cells, $\langle \hat{\Gamma} \rangle_6$, for WT and LatA treated cells. Apical tensions are not different in pentagonal and heptagonal cells, but basal tensions are higher (lower) at pentagonal (heptagonal) cells for WT cells. Treatment of LatA lowers this basal tension, as could be expected from the increased opening angles. The basal tension increase at pentagonal cells is negatively correlated with the relative angle increase at pentagonal cells (correlation $\rho_c = -0.29$ with p-value $p = 0.07$), supporting the notion of a complex balance between opening angles and basal tension. Linear regression on both the WT and LatA data suggests a relative angle increase of approximately 17% if the pentagonal basal tension is not different from the hexagonal one [Fig. 3(d)]. This finding indeed hints at an active compensation effect being responsible for the shape difference.

To explore potential mechanisms by which cells could adapt dynamically to their topological neighborhood, we determine the midplane areas as the ratio of cellular volume and height, V_{cell}/h [cf. Fig. 3(e)]. Similarly to Lewi's law in two-dimensional epithelia and foams, which states that area increases with the number of neighbors [44–46], we find approximately 10% smaller pentagonal and 15% larger heptagonal areas. This finding enables a direct feedback of tissue topology to the cell and could be linked to mechanosensing, similar to density sensing in intestinal epithelia [1,47]. Looking at the relationship of pentagonal basal tensions and areas, we find a negative correlation $\rho_c = -0.34$ (p-value $p = 0.03$): For smaller average pentagonal areas, we find increased basal tensions [Fig. 3(f)]. Linear regression suggests that this tension increase vanishes as areas become comparable to hexagonal areas, which is consistent with an area-driven feedback via mechanosensing facilitating a topology-dependent stabilization. Nonetheless, the underlying bulging instability is not as prominent in mouse colon organoids, implying the need for stabilization mechanisms also in our theoretical description.

D. Predictions of BVM for different shape stabilization mechanisms

Our experiments have revealed a coupling of tissue topology to cell shape and force generation. In particular, the data hint at an active basal stabilization mechanism that

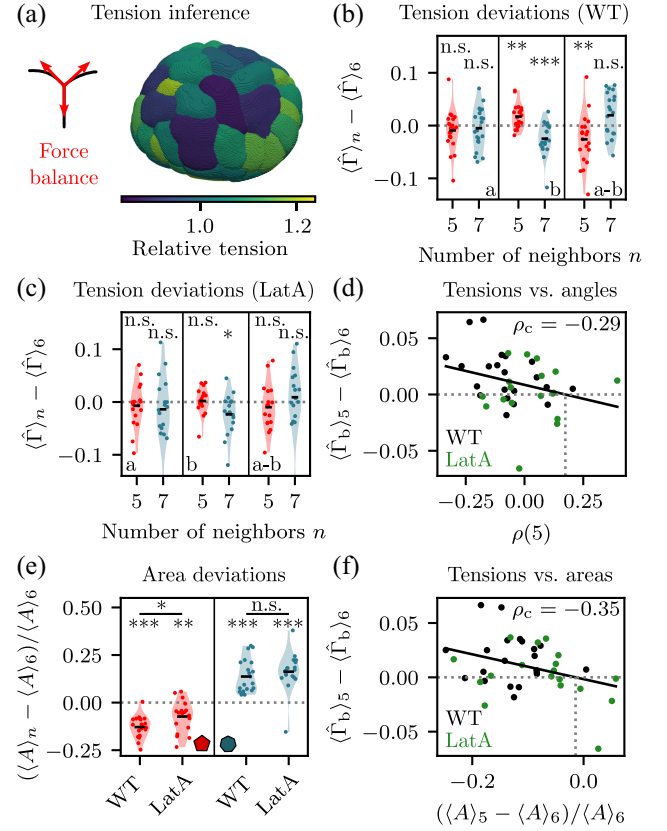


FIG. 3. Surface tension inference in mouse colon (mCol) organoids. (a) Relative tensions inferred from the dihedral angles of the interfaces, assuming force balance. For the reconstructed organoid, the outer (basal) tensions are shown. (b, c) Difference of the averaged inferred tensions $\hat{\Gamma}$ for n -gons from hexagons for apical (denoted in panels with a) and basal (denoted with b) interfaces for each WT (b) and LatA (c) organoid. The averaged cell-specific apicobasal tension difference is denoted as $a - b$ in the panels. (d) Pentagonal basal tension deviations versus relative angle increase of pentagons compared to hexagons $\rho(5)$ for WT and LatA samples. The solid line is the linear regression, which intersects a tension deviation of 0.0 at $\rho(5) \approx 0.17$. (e) Relative midplane area deviations of pentagons (red) and heptagons (blue) from hexagons for WT and LatA organoids. (f) Pentagonal basal tension deviations versus relative area deviations of pentagons for WT and LatA samples. The solid line is the linear regression, which intersects a tension deviation of 0.0 at $(\langle A \rangle_5 - \langle A \rangle_6) / \langle A \rangle_6 = -0.015$. The Pearson correlation coefficient in panels (d) and (f) is ρ_c . Stars indicate the following: * $p < 0.05$, ** $p < 0.01$, *** $p < 0.001$. Tests against 0 in panels (b), (c), and (e) are with the Wilcoxon test, and tests against each other in panel (e) are with the Mann-Whitney-U test.

controls cellular shapes at topological defects (pentagons) in organoids. This finding agrees with a very recent experimental report that basal contractility prevents extrusion in crypt regions [8], but it is challenging to identify the exact mechanisms given the complex situation in organoids. In order to better understand the physical processes that might be involved in cell extrusion and shape stabilization,

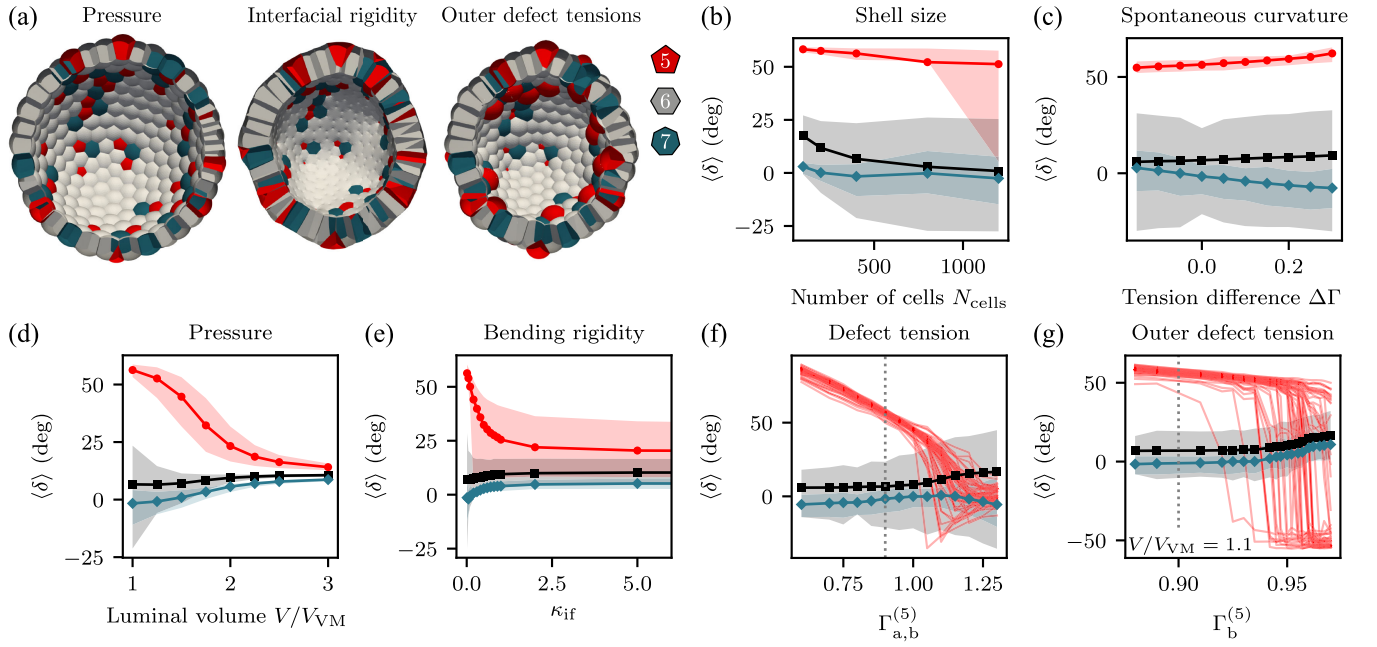


FIG. 4. Modulation of instability via different mechanisms explored through the BVM. (a) Shell from Fig. 1 simulated with additional effects: pressurized with $V/V_{\text{VM}} = 2.0$, rigidified with $\kappa_{\text{if}} = 0.7$, and pentagonal basal (outer) defect tension of $\Gamma_b^{(5)} = 0.957$. (b) Opening angles of pentagonal (red), hexagonal (black), and heptagonal cells (blue) as a function of the number of simulated cells. (c) Opening angles for different apicobasal tension differences $\Delta\Gamma = \Gamma_a - \Gamma_b$, corresponding to different spontaneous curvatures. (d) Opening angles for different luminal volume constraints, where $V/V_{\text{VM}} \approx 1$ corresponds to the case without a luminal volume constraint. (e) Opening angles for rigidified shells with different interfacial bending rigidities κ_{if} , where $\kappa_{\text{if}} = 0$ is the regular BVM. (f) Dependence on the basal apical tensions $\Gamma^{(5)} = \Gamma_a^{(5)} = \Gamma_b^{(5)}$ of pentagonal cells, with unchanged tensions in other cells. (g) Dependence on basal (outer) tension $\Gamma_b^{(5)}$ of pentagonal cells, which is changed compared to the outer tension of the other cells. All other tensions are unchanged compared to the reference. Markers are average angles, and shaded regions display 5% to 95% quantiles. Single red lines in panels (f) and (g) are individual pentagonal cells, and dotted vertical lines mark the reference case.

we turn back to the BVM. Figure 4 depicts the role of different tissue and cell properties on the opening angles at defects in simulated BVM shells. Figure 4(a) displays exemplary shell shapes in modified BVM shells. For increasing shell size, we find that the effect of the topological defects on the opening angles decreases [Fig. 4(b)]. In particular, for large shell sizes, we see that some pentagonal cells also display negative opening angles, which is effected by cells bulging into the lumen as a consequence of the reduced background curvature. Introducing spontaneous tissue curvature via an apicobasal tension difference $\Delta\Gamma = \Gamma_a - \Gamma_b$ modulates the opening angles slightly, but it cannot fully stabilize the cell shape [Fig. 4(c)], even for a spontaneous curvature radius of hexagonal cells that corresponds to the shell radius, here, $\Delta\Gamma \approx 0.14$. In the following, we therefore usually neglect spontaneous curvature, i.e., consider $\Gamma \equiv \Gamma_a = \Gamma_b$. Figure 4(d) studies a shell whose luminal volume V was prescribed to be a multiple of the luminal volume of the relaxed VM shell, V_{VM} , corresponding to an internal luminal pressure. Internal pressure is able to unfold pentagonal defect cells, cf. Fig. 4(a). Contrary to the different opening angles for different numbers of neighbors, cf. Fig. 1(c), increased pressure leads to convergence of the

average opening angles of pentagons, hexagons, and heptagons; see Fig. 4(d). In addition, we also see that the distribution of opening angles in pressurized BVM shells is less broad, more closely resembling the classical VM case. Thus, pressure can be used to suppress the buckling instability caused by the surface tensions. To estimate the experimental level of pressure-mediated stabilization, we calculate the theoretically expected aspect ratio [19,39] of hexagonal cells in the limit of flat interfaces (the classical VM) from inferred tensions [cf. Fig. 3(a); median relative tensions $\Gamma_a \approx 1.02$, $\Gamma_b \approx 1.06$, $\Gamma_l \approx 0.97$] to be $\psi_{\text{theo}} = h/A^{1/2} = 2^{1/2}3^{-1/4}(\Gamma_a + \Gamma_b)/\Gamma_l \approx 2.3$ with cell height h . Experimentally, we find $\psi_{\text{exp}} = h^{3/2}V_{\text{cell}}^{-1/2} \approx 1.3$, yielding an experimental to theoretical shape ratio of $\psi_{\text{exp}}/\psi_{\text{theo}} \approx 0.6$, which is consistent with a luminal pressure stretching the epithelium. Comparing this ratio to the pressurized theoretical BVM case, we find similar values for a luminal volume of $V/V_{\text{VM}} \approx 2.25$ ($\psi_{\text{exp}}/\psi_{\text{theo}} = 0.62$), where cells have not fully unfolded. If we take average surface tensions of intestinal organoids from the literature, which are of the order $\langle \Gamma \rangle \approx 5 \text{ nN}\mu\text{m}^{-1}$ [48], we can roughly estimate the luminal pressure in the limit of flat interfaces (the classical

VM) to be $P \approx 440\text{Pa}$ (see Appendix D for details). In our case, these pressures alone cannot fully explain smaller opening angles of pentagonal cells as observed experimentally.

To highlight the role of interface curvature, we added an interfacial bending energy to the BVM minimization for each interface Ω_{if} ,

$$E_{if} = \frac{\kappa_{if}}{2} \int_{\Omega_{if}} H_{if}^2 dS, \quad (3)$$

with interfacial bending rigidity κ_{if} and (total) mean curvature of the interface H_{if} . This model penalizes interfacial curvature energetically. Figure 4(e) shows the opening angles, whose values quickly approach each other as κ_{if} is increased. However, rigidifying the interfaces does not achieve as much convergence of the opening angles toward each other as luminal pressure. In simulations, we find the cells to have less curved interfaces, interpolating between the VM and the BVM, but this outcome does not agree with experimental observations, cf. Fig. 2.

Motivated by our experiments, we further change the apical (inner) and basal (outer) tensions of the pentagonal defect cells to be different from the other cells to introduce topology-dependent variability [Fig. 4(f)]. In this case, we see that increasing the defect tension leads to a reduction of pentagonal opening angles, where a transition through a shape consistent with hexagons toward negative opening angles occurs. Note that reducing the tension can drastically increase the pentagonal opening angle, already hinting at a potential extrusion mechanism through a tension decrease.

Figure 4(g) shows the effect of increased outer (basal) tensions at pentagonal cells, as found experimentally in the mCol organoids. We also assume a nonvanishing luminal pressure, as this is likely to reflect the experimental situation. We find that the cells transition through a state consistent with hexagonal cells toward a state with approximately opposite opening angle, implying full cell shape control, cf. shell in Fig. 4(a). Note that the transition point is cell specific, as it likely depends on the details of the neighborhood topology. This feature suggests a potential feedback mechanism between local topology (and induced stresses) and interfacial tensions.

Beyond surface tensions, we also consider an increase in global basal line tensions and find that it, too, can stabilize cell shapes; see Appendix E. The basal stabilization mechanism observed in experiments is thus proven to enable reliable control over the cell shape.

E. Shift of icosahedral buckling threshold

The collapse of the luminal interface at pentagonal cells is reminiscent of the collapse observed earlier in the buckled icosahedral case in the VM for large apicobasal tensions Γ and thus cell heights [19]. In this case, a tissue-scale

azimuthal stress from the pentagonal defect configuration is relaxed through conical buckling for large tissues, an effect described before for elastic shells, in general [23], leading to an icosahedral shape due to this effect occurring at the 12 pentagonal defects in the hexagonal tiling of the sphere. This system-scale instability is well known from viral capsids [20,21] and buckyballs [22]. On a single-cell scale, at these topological defects, luminal interfaces can collapse due to the large necessary cellular opening angle at the icosahedral tips, quite similarly to what we observe in the BVM.

To be able to compare with these results for the VM, we now also consider spherical shells with icosahedral symmetry, following the Casper-Klug construction [49]. The Caspar-Klug indices (T_1, T_2) describe the steps in the underlying hexagonal lattice to go from one pentagonal defect cell to another with respect to lattice vectors with an opening angle of $\pi/3$. If we compare the results for the VM [Fig. 5(a)] versus the results for the BVM [Fig. 5(b)], we see that the buckling instability in the BVM is much stronger for the same number of cells. It also occurs at smaller radii than the buckling instability in the VM, which only occurs for large system size. To quantify this effect, we measure the asphericity

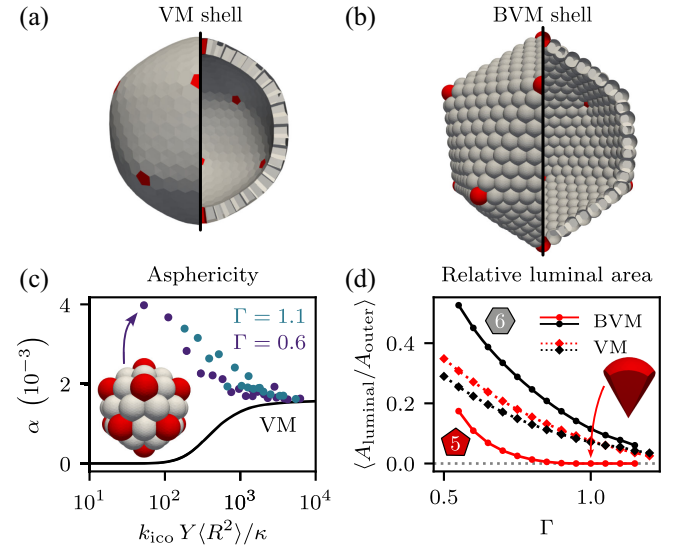


FIG. 5. Lowered buckling threshold for icosahedral instability in the BVM compared to the VM. (a, b) Icosahedral shells with Caspar Klug indices (8, 0) with tension $\Gamma = 0.6$ in the VM (a) and the BVM (b). (c) Asphericity α in the BVM shell (symbols) and of the continuum curve describing the VM (taken from Refs. [19,20]). The inset shows a shell with size (2, 0), $\Gamma = 0.6$. Radii are rescaled by the Föppl-von Kármán number and with a correction k_{ico} from nonlinear elasticity. (d) For a small shell [Caspar Klug indices (2, 0)], the average ratio of luminal and outer interfacial area in the BVM (solid) and VM (dashed). The inset shows a pentagonal cell with collapsed luminal interface ($\Gamma = 1.0$). Pentagonal cells are shown in red.

$$\alpha = \frac{\langle (R - \langle R \rangle)^2 \rangle}{\langle R^2 \rangle}, \quad (4)$$

where R is the radial distance of a cell (determined as the average of the centers of mass of the apical and basal interfaces) from the shell center and $\langle \cdot \rangle$ is the multicellular average. For the VM, this quantity has been shown to follow a continuous curve as a function of the rescaled shell size, corresponding to the Föppl–von Kármán number $\gamma_{\text{FvK}} \equiv Y_{\text{eff}} R^2 / \kappa$, with bending rigidity κ , radius R , and (effective) Young's modulus Y_{eff} , corrected here for effects from nonlinear elasticity with a numerical factor k_{ico} [19]. As shown in Fig. 5(c), the asphericity is much larger for small shells in the BVM compared to the VM. Details on the axis scaling are given in Appendix H.

We also find that the asphericities in the BVM do not fall on one curve but rather deviate from an average curve, which results from both the broad distributions of hexagonal opening angles described above, and therefore less regular structures, and from cellular positioning for small cell numbers on the icosahedral surface due to under-sampling of the icosahedral surface. Similarly to the VM, we have a collapse of the luminal interface for larger Γ , which is much more pronounced in the BVM with a stronger deviation of the luminal-to-outer area ratios between pentagonal and hexagonal cells for small spheres, cf. Fig. 5(d). Even shells as small as (2, 0), i.e., only one hexagonal cell between the pentagonal cells, show a big difference between pentagonal and hexagonal cells and thus adopt a strongly icosahedral shape. We thus see that the buckling threshold is reduced to effectively one neighboring cell in the BVM, rendering the icosahedral tissue-scale instability much more relevant in the case of curved cellular interfaces in the BVM.

III. CONTINUUM THEORY FOR EXTRUSION OF A DEFECT CELL IN A MEAN-FIELD TISSUE

A. Parametrizing the defect cell

To understand the nature of the significant lowering of the icosahedral buckling threshold for the BVM compared to the VM and expand the model to capture cell extrusion, we now investigate the BVM analytically in a continuum setting. We focus on the icosahedral tip and the conical instability to first develop the continuum model in a setting amenable to comparison to numerical simulations. At the site of the disclination defect, we excise the continuous material and directly consider the energy of the pentagonal cell, following the description of the BVM, cf. Eq. (1). Figure 6(a) depicts the geometry that we assume.

At the center of the disc of material is the defect cell, which we consider first. We assume rotational symmetry, and thus the cutout of the tissue is the base of a cone. The cell is allowed to bulge outward via interfacial curvature,

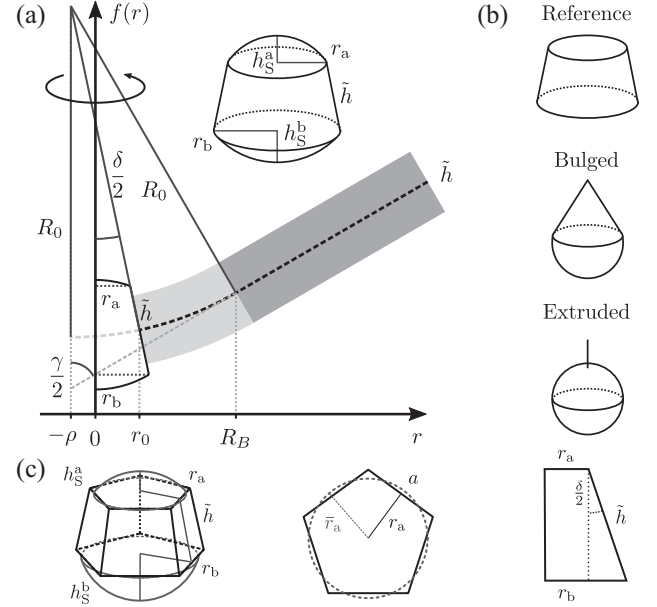


FIG. 6. Schematic depiction of the defect cell in a mean-field tissue. (a) Tissue profile as a radial function in cylindrical coordinates. At the center is the defect cell, which consists of the base of a cone with spherical caps on the apical and basal sides. This defect is surrounded by mean-field tissue that forms a spherical cap segment (light gray) and, for larger distance, a cone (dark gray). (b) Model described by different defect cell geometries, namely, the vertex-model reference case ($h_S^a = h_S^b = 0$), a completely bulged case ($r_a = 0$, $\delta/2 > 0$), and a fully extruded case ($r_a = 0$, $\delta/2 = 0$). (c) Geometrical correction for pentagonal cells, where we assume a pentagonal pyramid base with two spherical caps. The radii of the caps, $\tilde{r}_{a,b}$, are chosen such that the area of the caps' disc bases, the circles, match the pentagonal areas (center), and the slant height $\tilde{h}$ is considered at the lateral faces' centers.

which is captured by considering spherical caps at the apical and basal sides of the defect cell.

With the chosen parametrization, we are able to describe different defect cell morphologies, i.e., a vertex model case with flat apical and basal interfaces, a bulged case with a collapsed interface, and an extruded case [Fig. 6(b)]. Because of the pentagonal nature of the defect, we modify the conical shape to reflect the pentagonal structure by considering a pentagonal pyramid base; see Fig. 6(c). Thus, we introduce a geometrical factor ζ_{geom} , which relates the basal and apical polygonal edge lengths, b and a , respectively, to the areas such that, for the apical area, $A_a = \zeta_{\text{geom}} a^2$ and similarly for the basal side. Additionally, we introduce η_{geom} , which relates the radius of a regular polygon's incircle $r_{a,b}$, i.e., the largest circle that fits into the regular polygon, to the edge length, e.g., for the apical side $r_a = \eta_{\text{geom}} a$.

For a bulged pentagonal cell, we consider the pyramid base with apical (or basal) spherical caps. We assume the caps have the same disc area as the apical-basal pentagons,

neglecting the slight mismatch of the two different shapes; cf. Fig. 6(c). This case is consistent with simulated cell shapes in the BVM as shown in Fig. 1(a). For the spherical cap, we denote the base disc radius with t_a (or t_b for basal); the cap's base area is πt_a^2 , which implies $\pi t_a^2 = \zeta_{\text{geom}} a^2$ for the apical side (and analogously for the basal side). The radius of the apical spherical cap's base then relates to the edge length via $t_a = \sqrt{(\zeta_{\text{geom}}/\pi)}a$ (analogously for the basal side).

Henceforth, we denote the side to which the tissue buckles (positive z) the apical side and the opposite side the basal one, as this corresponds to the situation in the experimental mCol organoids when considering closed shells. Note that the sign of curvature at the villus tip in the small intestine is the opposite. The opening angle of the pyramid at the lateral face centers, $\delta/2$, the apical inradius r_a , and the tissue height $\tilde{h}$ define the base part. To consider the energy change due to extrusion and bulging, we define, as the reference case, the vertex model case, where the defect cell is not bulged and induces the same curvature as the tissue would. The spherical caps of the defect cell are uniquely defined by the basal and apical radii and the height of the caps $h_S^{a,b}$ (cf. Fig. 6).

The radius of curvature of the spherical caps $R_S^{a,b}$ depends on the pressure differences between the cell and the surrounding medium $\Delta P_{a,b}$ at the corresponding side via the Young-Laplace law,

$$\Delta P_{a,b} = \frac{2\Gamma_{a,b}}{R_S^{a,b}}. \quad (5)$$

If we assume equal pressure differences, i.e., no pressure difference between the basal and apical medium, the radii are related via

$$\frac{R_S^a}{\Gamma_a} = \frac{R_S^b}{\Gamma_b}. \quad (6)$$

As described earlier, we assume $\Gamma_a = \Gamma_b$ in the analytical treatment, i.e., no apicobasal polarity. Using this relation, we express the basal cap height via the apical one, assuming that extrusion only occurs on the outer (basal) side (similar to the sign of curvature at the intestinal villus tip) and therefore choosing the smaller of two possible cap heights.

The volume of the cell then implicitly defines h_S^a via

$$1 = \frac{\tilde{h} \cos(\delta/2)}{3} \zeta_{\text{geom}} (a^2 + ab + b^2) + \frac{h_S^b}{6} (3\zeta_{\text{geom}} b^2 + \pi h_S^{b2}) + \frac{h_S^a}{6} (3\zeta_{\text{geom}} a^2 + \pi h_S^{a2}), \quad (7)$$

with the lateral (slant fraction) height $\tilde{h}$. The last two terms correspond to the spherical caps and vanish for the

nonbulged reference cell as $h_S^b = h_S^a = 0$. The apical and basal sides are related to each other via the opening angle, $b = a + \sin(\delta/2)\tilde{h}/\eta_{\text{geom}}$; cf. Fig. 6(c).

In a continuum approach, we would not resolve the defect cell, and the surrounding spherical cap with spherical radius R_0 would determine the opening angle, i.e., $\sin(\delta/2) = \eta_{\text{geom}} \{1 + [(R_0 + \tilde{h}/2)/(R_0 - \tilde{h}/2)]\} a/2R_0$. However, resolving the defect cell explicitly allows us to consider a vertex model later, which goes beyond this approximation, even if we do not allow for interfacial curvature.

With known cap heights, we can write down the energy of the bubbly defect cell by adding the lateral surface energies and the surface energies of the spherical caps to obtain

$$E_{\text{core}}^{\text{bulged}} = \Gamma_a (\zeta_{\text{geom}} a^2 + \pi h_S^{a2}) + \Gamma_b (\zeta_{\text{geom}} b^2 + \pi h_S^{b2}) + \frac{n\tilde{h}}{4} \left(2a + \frac{\tilde{h}}{\eta_{\text{geom}}} \sin(\delta/2) \right), \quad (8)$$

where the cell is given by a regular n -gon. For a pentagonal defect, we have $n = 5$, $\zeta_{\text{geom}} = \sqrt{5(5 + 2\sqrt{5})}/4$, and $\eta_{\text{geom}} = 1/(2\sqrt{5 - \sqrt{20}})$. We assume that the lateral slant height of the pyramid base is the equilibrium height of the surrounding tissue $\tilde{h}$, and thus the only two free parameters are a and δ . To describe the full tissue, we additionally need the background curvature R_0 , which relates to the buckling radius R_B via the conical opening angle γ ; cf. Fig. 6(a).

Figure 7 shows the dimensionless core energy, Eq. (8), for different opening angles $\delta/2$ as a function of the apical edge length a for (a) $\Gamma = \Gamma_a = \Gamma_b = 0.65$ and (b) 1.0, respectively. Edge length a is bounded from above by the volume constraint in the unbulged vertex model case with $h_S^b = h_S^a = 0$. For each angle, the energy initially decreases if we decrease a with concomitant outward bulging (top panel), as the area-to-volume ratio decreases. For even smaller a , the energy usually increases again, suggesting that a configuration with curved interfaces is always energetically more favorable.

Increasing the angle leads to an increase in the core energy in the vertex model case (cf. dashed line at which the curves end in the top panel) and shifts the minimum of the energy toward smaller a if we allow for bulging. If we compare the vertex model energy with the minimal energy in the bubbly vertex model for given opening angles (and thus surrounding tissue curvature), we find a smaller energy in the latter case; see lower panel.

Strikingly, as we increase $\delta/2$ and thus bend the defect cell, we either obtain a smaller increase of energy in the bubbly case compared to the vertex model ($\Gamma = 1.0$) or even an energy release ($\Gamma = 0.65$). As this result is related to the bending rigidity, it drastically changes the behavior of the overall tissue. The potential to bulge outward

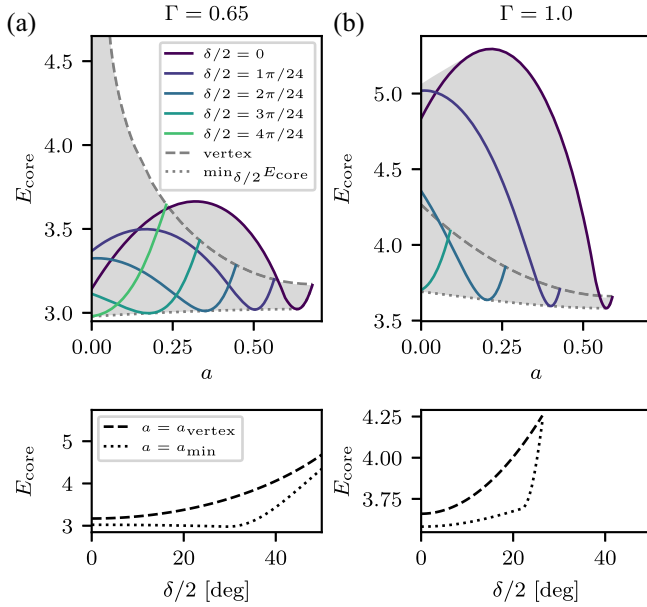


FIG. 7. Core energy for the different defect cell configurations for (a) small apical-basal surface tension Γ and (b) large Γ . The solid colored curves in the upper panels are the energies as a function of apical edge length a for fixed opening angles $\delta/2$. The curves end at the vertex line (dashed) at which a is maximal for the angle such that the apical and basal faces do not bulge. All possible states and energies that can be reached with $\delta/2$ and a such that the volume constraint is fulfilled are shaded in gray. The minimum line (dotted) indicates the minimal energy for a given a . In the lower panels, the vertex line and the minimum line are parametrized with $\delta/2$. The minimum line is obtained by finding the edge length a_{min} that minimizes E_{core} for a given $\delta/2$. We consider a pentagonal defect with equilibrium tissue height of the hexagonal vertex model, i.e., $\tilde{h} = 2^{1/3}3^{-1/6}(2\Gamma)^{2/3}$.

therefore decreases the necessary bending energy in the defect core and can even lead to an energy release. This energy can then compensate for energy contributions in the surrounding tissue, allowing for a different equilibrium shape than in the two-dimensional elastic sheet. This effect explains the observed concentration of Gaussian curvature in the defect cell, as it is energetically cheaper to go into the buckled state for an icosahedral shell.

B. Deformation of the mean-field tissue

For the surrounding mean-field tissue, we assume that the inner radius of the inner tissue segment [cf. light gray part in Fig. 6(a)] corresponds to the polygonal inradius of the defect cell, neglecting mismatch contributions stemming from the shape difference between the cellular pentagonal midplane and circular cutout. This inner cap segment is assumed to follow a circular arc with radius of curvature R_0 in cylindrical coordinates, cf. Fig. 6(a). An outer segment (cf. dark gray part) is assumed to be conical, in agreement with the result in moderately bent plates with a disclination defect.

We use the description of a moderately bent plate as derived in Ref. [19]. The energy change due to the defect deformation can then be split into different contributions. First, the bending energy will change, both for the mean and Gaussian curvature contributions, due to changes in the curvature radii and opening angles in the inner cap segment. Second, the stretching energy will change as the azimuthal strain stemming from the pentagonal defect can be relaxed through changes in the curvature. For simplicity, we neglect radial stretching contributions at the interfaces by considering stress-free boundary conditions at the contact points of the defect cell, cap segment, and outer cone segment, and at the outer edge of the cone. This approach is justified numerically in Appendix F.

Introducing the bending energy density e_{bending} and the stretching energy density e_{stretch} , the total energy of the tissue Ω then reads

$$E = \int_{\Omega} (e_{\text{bending}} + e_{\text{stretch}}) dS. \quad (9)$$

For the spherical topology of the tissue that we examine here, we consider the conical deformation of a disc with the pentagonal defect in the center, as this has been shown to capture the icosahedral instability of such a regular sphere [19,20].

The different energy contributions for bending and stretching for this mean-field tissue parametrization are derived in Appendix F. This parametrization with a mean-field model has been verified on icosahedral vertex model shells without interfacial curvatures, where the mentioned decomposition into a defect cell and a mean-field tissue is able to describe the tip much better than an approach purely based on continuum theory; cf. Appendix G.

C. Extrusion of a bubbly defect cell in a mean-field vertex model tissue at icosahedral tips

We now consider a bubbly defect cell in a mean-field tissue of icosahedral symmetry; i.e., h_s^b and h_s^a may deviate from zero.

In the following, we assume the mean-field tissue parameters to be from the classical (nonbubbly) VM, thus only allowing for bulging of the defect cell. In principle, however, mean-field parameters could be derived for the BVM in a similar fashion as for the VM, or they could be taken from experimental measurements for the tissue properties.

The total energy is the sum of the contributions from the defect, the cap, and the outer cone region. We consider the energy difference from a reference state to quantify the energy release due to bulging by considering the energy difference ΔE from the vertex model case with minimal energy. The latter corresponds to the full defect model described in Appendix G; see Fig. 12.

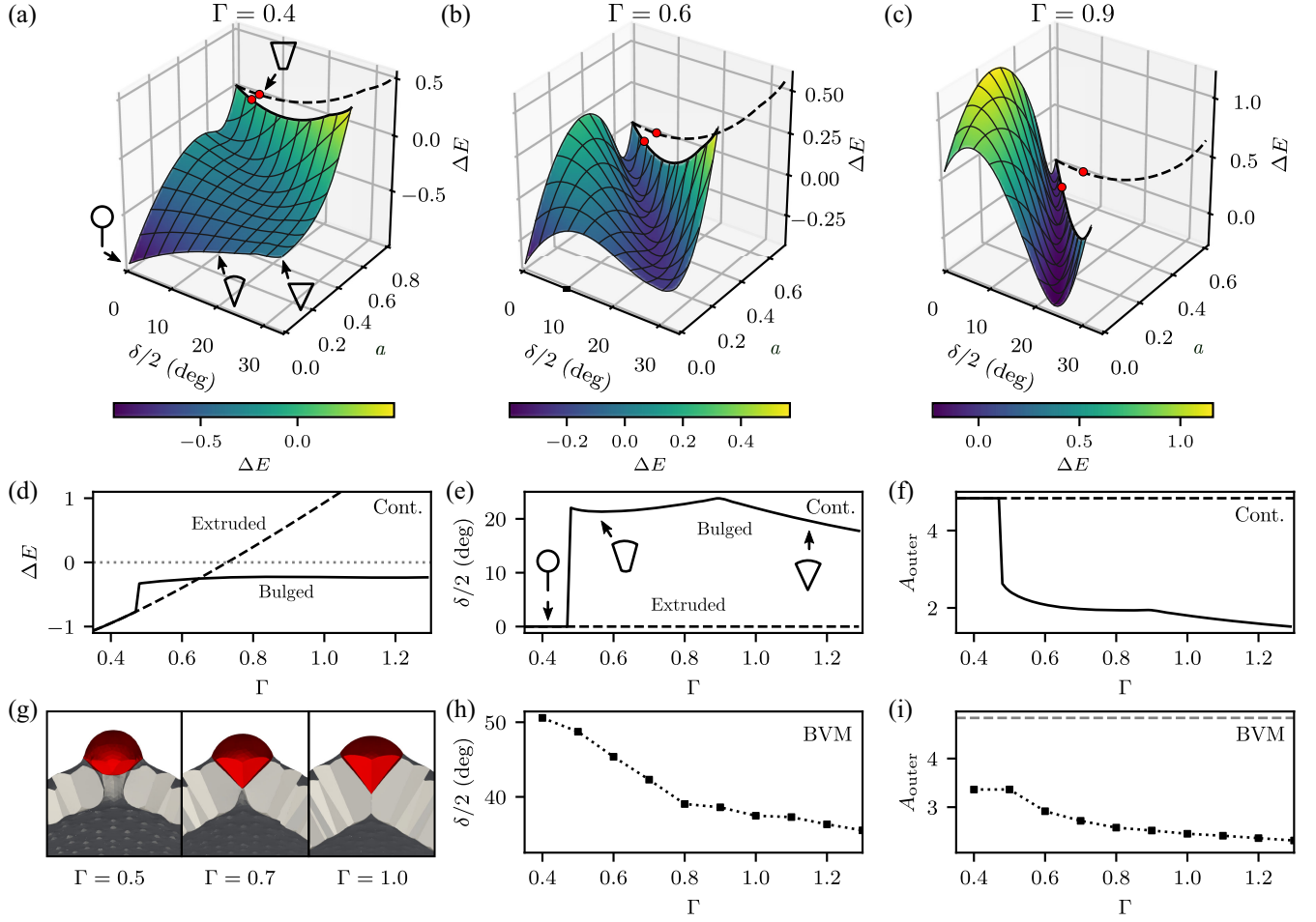


FIG. 8. Energy landscapes in configuration space and the resulting defect cell shapes. (a)–(c) Energy landscapes with coordinates $(\delta/2, b)$ for different surface tensions Γ . The black line marks the vertex model limit without bulging of the defect cell, and the red circle marks the minimum in this limit. The dashed line is the projection onto the $(\Delta E, \delta/2)$ plane of this limit. The color-coding is for the energy difference ΔE with respect to the reference energy, which is the vertex model minimum. The corners of the landscape at $(\delta/2, b) = (0, 0)$ and $(\delta_{\text{max}}/2, 0)$ correspond to the extruded and maximally bent situations, respectively; cf. Fig. 6. (d) Continuum model energy differences from the VM reference ΔE of the minimal energy configuration (solid), starting from the VM reference, and of the extruded configuration (dashed). (e), (f) Corresponding defect opening angles $\delta/2$ and outer interfacial area A_{outer} . (g) BVM with interfacial bending rigidity, where the defect cell is extruded for small apicobasal tension $\Gamma = 0.5$ and partially extruded for large tensions $\Gamma = 0.7$ and $\Gamma = 1.0$. (h), (i) Corresponding opening angles and outer interfacial area for the BVM. Shells of Caspar-Klug size $(T_1, T_2) = (12, 0)$ are shown with bending rigidity $\kappa_{\text{if}} = 5$ in dimensionless units.

We find that, for all configurations characterized by $\delta/2$ and a , the energy has a unique minimum in the tissue buckling radius R_B . We assume this radius to be minimal for every defect cell shape and thus can describe all possible shapes by the variables $(\delta/2, a)$, which constitutes the configuration space. Both parameters are bounded from below by 0, where $(0, 0)$ is possible and corresponds to the fully extruded case, and from above by the volume constraint of the defect cell.

Figures 8(a)–8(c) show energy landscapes in configuration space for different surface tensions Γ . For now, we assume identical cell properties; i.e., the tension in the defect cell matches the one of the surrounding tissue. We notice that, by bulging outward at the defect, the energy of

the tissue can be reduced considerably in comparison to the classical vertex model reference case. For small Γ , more energy can be released by bulging outward than for larger Γ . For $\Gamma > 0.5$, we observe that two minima form in the energy, one in the extruded case $(\delta/2, a) = (0, 0)$ and one in a bulged case with very small a . The energy minimum in extrusion is higher than the bulged one for larger Γ , and an energy barrier has to be overcome from both the reference and the global minimum configuration to extrude the defect.

This effect is to be expected because we have assumed that the bulged membrane has surface tension Γ , and thus the fully extruded cell (surrounded by said membrane) will have a higher energy due to this larger surface tension. For

$\Gamma < 0.5$, on the other hand, we find that the extruded case is minimal in energy and can directly be reached without passing over an energy barrier. The case $\Gamma = 0.5$ is special because it corresponds to cells with an aspect ratio of approximately 1, as the lateral energies only contribute half to the total energy [cf. Eq. (1)]; thus, energy costs for lateral and apicobasal interfaces do not differ. For $\Gamma < 0.5$, the lateral membrane is effectively more expensive than the apical membrane.

We now numerically determine the (local) minimal energy configurations for different Γ , starting from the VM reference state. Figures 8(d)–8(f) depict the resulting minimal energy configurations. Above a critical tension Γ , the luminal (basal) interface collapses but the angle remains finite, indicating a completely bulged state. The fully extruded state at $(\delta/2, a) = (0, 0)$ always represents a possible configuration, but its energy increases with Γ . Looking at the energy difference from the VM state, a configuration with interfacial curvature is always energetically more favorable; cf. Figs. 8(e) and 8(f). This finding is not surprising as the surface area can be minimized by the cell becoming more spherical. The extruded case, on the other hand, requires more energy for large apicobasal tension, as we need to wrap the cell in the energetically (relatively) expensive apical interface. Lowering the tension allows for an energy release in the extruded case, which can become lower than in the bulged state. Note, however, that a bulged state persists beyond the point at which the extruded case is energetically more favorable and that an energy barrier exists in configuration space, which prevents the cell from being directly extruded from the reference VM configuration; cf. Fig. 8(b).

Considering icosahedral shells in the BVM and adding a bending rigidity across triangulation points of hexagonal cells' interfaces enables a numerical simulation similar to our continuum approach of a bubbly defect cell surrounded by VM cells; cf. Fig. 4(g). We thus consider an additional bending energy, Eq. (3).

Figure 8(g) shows configurations of the defect cell in this BVM with bending rigidity for different apicobasal tensions Γ . For $\Gamma \leq 0.5$, we see that, indeed, the pentagonal cell is pushed out of the monolayer, only being attached because the lateral interfaces cannot be broken in the simulations. This finding is also visible in the parameters quantifying the configurations in Figs. 8(h) and 8(i). The luminal interfaces bend outward despite the energetic penalty to facilitate this extrusion process.

For intermediate Γ , the neighbor cells are deformed due to finite κ_{if} , while for large Γ , a completely bulged state is found, as predicted by the continuum model.

D. Extrusion at pentagonal and hexagonal defects in spherical shells with curvature screening

We generalize our continuum description to spherical shells with curvature screening, as we have also seen the

effect of topological defects in less structured monolayers; cf. Fig. 1. Thus, we assume the topological defect to only induce azimuthal strain locally around the defect. In simulations, we see that pentagons with one neighboring ring of hexagonal cells exist in spherical shells due to the necessary excess topological charge (12 more pentagons than heptagons); cf. Fig. 1. Thus, we consider a defect cell, surrounded by a cap region with varying radius of curvature but fixed size (i.e., fixed buckling radius R_B), which in turn is surrounded by background tissue as depicted in Fig. 9(a). The background tissue is assumed to be mechanically relaxed and not considered in the following: The curvature is screened accordingly so that a spherical configuration is obtained [cf. Fig. 1(b)], and topological defects at most affect their neighbors. The inradius for a hexagonal cell is

$$r_0^{\text{hex}} = \sqrt{\frac{1}{\bar{h}\zeta_{\text{geom}}^{\text{tissue}}}} \eta_{\text{geom}}^{\text{tissue}} = 3^{-1/6} 2^{-2/3} (2\Gamma)^{-1/3}. \quad (10)$$

For the radius of the cap along the radial axis, we now assume $R_B = 3r_0^{\text{hex}}$, i.e., one neighboring ring of cells. The background curvature radius R_{bg} only enters via the opening angle (similarly to the conical opening angle), i.e.,

$$\pi/2 - \gamma/2 = \arcsin(R_B/R_{\text{bg}}). \quad (11)$$

For this configuration, the total energy is now identical to the contributions described in Appendix G [cf. Eq. (G2)] but without the cone energy and with the prescribed R_B and new angle $\gamma/2$, i.e.,

$$E_{\text{tot}} = E_{\text{core}} + E_{\text{bending}}^{\text{mean}} + E_{\text{bending}}^{\text{Gauss}} + E_{\text{stretch}} + E_{\text{bdry}}. \quad (12)$$

We now also want to consider extrusion of a hexagonal cell, for which we have no azimuthal stretching [a disclinity of $s = 0$ in Eq. (F12) in Appendix F] and consider $n = 6$ lateral faces. The geometric constants then read $\eta_{\text{geom}} = 3^{1/2}/2$ and $\zeta_{\text{geom}} = 3^{3/2}/2$. We have found that lower apicobasal tensions lead to extrusion because the cell has to be completely engulfed by the apical interface in the extruded case. This finding suggests that one potential mechanism of extrusion is the lowering of apicobasal tension in the defect cell. Figure 9(b) shows the energy differences from the VM reference for different defect tensions Γ_{defect} in a mean-field tissue of tension $\Gamma = 0.9$. We see that pentagonal energies are generally smaller in deformed states due to the relaxation of azimuthal strain when localizing Gaussian curvature in the defect.

For both hexagons and pentagons, we see that the extruded and bulged states become energetically more favorable as Γ_{defect} decreases. Below a critical tension

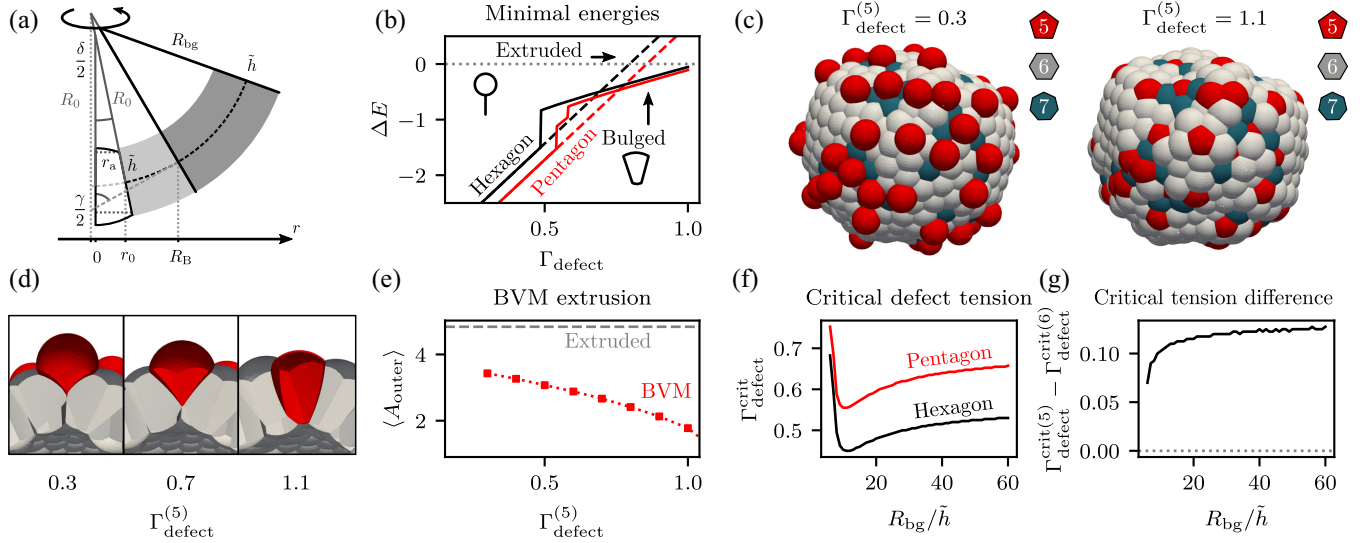


FIG. 9. Extrusion through lower defect tension in hexagons and pentagons. (a) Schematic depiction of the defect in mean-field tissue with background curvature (radius R_{bg}). (b) Energy difference for continuum model configurations with different defect tensions Γ_{defect} from the VM case with identical defect tension $\Gamma_{\text{defect}} = \Gamma$. The solid line is the local minimization from the VM case; the dashed line is the extruded case. Red is the pentagonal and black the hexagonal defect cell. (c) BVM shells from Fig. 1 in which pentagonal cells have different apicobasal tensions $\Gamma_{\text{defect}}^{(5)}$, displaying partially extruded morphologies that depend on the tension. Pentagons are shown in red, heptagons in blue. (d) Cross section of a pentagonal cell surrounded by hexagonal cells being extruded upon lowering of $\Gamma_{\text{defect}}^{(5)}$ in the BVM. (e) Average outer interfacial area of BVM shells. The horizontal dashed line corresponds to fully extruded cells. (f) Change of the critical defect tension $\Gamma_{\text{defect}}^{\text{crit}}$ in the continuum model, at which the bulged state becomes unstable toward extrusion for different background radii R_{bg} . (g) Difference of critical tensions in pentagonal and hexagonal cells in panel (f). The parameters are $\Gamma = 0.9$; (b), (f) $R_{bg}/\tilde{h} = 5$.

$\Gamma_{\text{defect}}^{\text{crit}}$, the bulged state becomes unstable toward the extruded state.

In BVM simulations, we find strong outward bulging reminiscent of intermediate states in cell extrusion upon lowering of the apical and basal defect tension, without achieving full extrusion due to the inability of closing the hole forming in the simulations [Figs. 9(c)–9(e)]. In Fig. 9(f), these critical tensions for pentagons and hexagons are plotted as functions of the background curvature radius $R_{bg}/\tilde{h}$. For a small radius, the critical tension is high, implying that extrusion is easy. The reason for this is that the volume constraint does not allow for a full relaxation of the Gaussian curvature in the defect because we have assumed that $\delta/2 < \pi/2 - \gamma/2$; i.e., the cap region cannot accommodate negative curvature. This feature leads to a less stable bulged state and thus earlier instability toward extrusion. For intermediate background radii, the bulged state is able to stabilize by relaxing the remaining Gaussian curvature in the defect. Increasing the background radius further then leads to less stability of the bulged state because the possible opening angle is small.

In the pentagonal cell, we see an intermediate state in Fig. 9(b): For intermediate tensions, a local minimum of full bulging ($a = 0$, $\delta/2 > 0$) forms, which then becomes unstable toward the extruded case for lower tensions [cf. Figs. 8(a) and 8(b)]. In general, we see that, for decreasing

defect tensions in both pentagons and hexagons, the energy landscape loses local minima corresponding to bulged states, as the extruded case becomes increasingly more energetically favorable.

Figure 9(g) depicts the difference in the critical defect tensions of pentagons and hexagons, $\Gamma_{\text{defect}}^{\text{crit}(5)} - \Gamma_{\text{defect}}^{\text{crit}(6)}$. The necessary tension decrease for overcoming the energy barrier for extrusion in hexagons is approximately 10% of the total tension higher than for pentagonal cells. While this suggests pentagonal cells are more prone to extrusion, it does not constitute a difference large enough for other cellular processes to not be able to overcome the barrier for hexagons as well. Introducing additional luminal line tensions yields a destabilization of a bulged state with a finite-size luminal interface; however, it does not contribute to the extrusion process without the involvement of lateral tensions; see Appendix E.

IV. DISCUSSION

Here, we have introduced the bubbly vertex model for multicellular epithelial monolayers as a generalization of the standard vertex model. In contrast to the VM, the BVM allows for curved interfaces, which are known to occur in flat epithelia and even more so in curved ones. Because curved epithelial sheets naturally have topological defects

due to Euler's polyhedron theorem, we studied if cell bulging and extrusion might occur naturally at such defects. By comparing the BVM with the VM, we found that allowing for cell interface curvature dramatically lowers the conical buckling threshold around such defects and yields a bulging instability.

To explore the significance of these theoretical predictions, we have performed experiments with spherical mouse colon organoids, which are an appropriate model system for the spherical epithelia described by our theory. Because extrusions are hard to capture due to their speed, we focused on cell shape at topological defects. We found that, indeed, cell shape is different at pentagonal cells, but differently from the predictions of the BVM, the opening angles were decreased rather than increased, which suggests the existence of some mechanisms for stabilization. Very recently, basal contractility has been reported as such a mechanism [8]. Leveraging a force inference pipeline that we have developed earlier [39], we indeed found an increased basal interfacial tension. Perturbation of the actomyosin cytoskeleton increased the bulging, showing its involvement in actively stabilizing cell shape. We found pentagonal midplane areas to be smaller and negatively correlated with basal tensions, potentially facilitating stabilization via mechanosensing. In simulations, we found this strategy of cell-specific tension variability to yield full control over cell shapes. We also showed that tissue-scale forces must be at play, in particular, pressure, which we estimate to be of the order of 440 Pa. This estimate is 1 order of magnitude smaller than reported for epithelia compressing hydrogel spheres [50] but likely larger than in mouse small intestinal organoids, which use lower pressures as a mechanism for budding [48].

In order to better understand the physical mechanism behind the instability, we developed a theory of a discrete defect cell in a mean-field tissue. This approach provides a mechanistic model for extrusionlike cellular shapes; in addition, it suggests that lowering the apicobasal tension in the cell to be extruded, and thus lowering the contractile forces and the mechanical stability of its cortex, provides a mechanism to extrude cells. This finding is consistent with recent experiments, showing extrusion upon cortical ablation in intestinal organoids and tension-related extrusion in intestinal villi [8]. Line tensions, which could arise from supracellular actin structures, facilitate extrusion but do not allow for the extrusion of partially extruded or fully bulged cells as they are observed in some experimental cases.

Our model makes experimental predictions concerning a facilitated extrusion at pentagonal cells and a curvature dependence of the necessary tension decrease for extrusion. These predictions could be tested in the future in systems such as intestinal organoids in scaffolds [51], where the intestinal epithelium is reconstituted by considering an intestinal monolayer on a 3D substrate with similar shape as found in the intestine. In particular, in this approach, cell

extrusion has indeed been found to occur at the villus tips [51]. It has recently been used to show that tissue tension and cytoskeletal stability play a role in extrusion [8]; in the future, it could be used to investigate our proposed dependence on the radius of background curvature. Recently, two-dimensional systems have become available to study extrusion [52], and instabilities could be investigated in such a setting as well. Intestinal organoids usually display an inverted sign of curvature compared to the villus domain in the intestine with the apical side pointing toward the lumen. Polarity-reversed organoids, however, have the opposite apicobasal orientation and thus the correct curvature. Therefore, they constitute another possible test system to study connections of single cell shape and tissue curvature [53,54]. Although hard to implement in practice, live cell microscopy might be used to capture extrusion events in more detail and to correlate them with the topological environment. It would be interesting to also extend our theoretical framework to include more dynamic processes, such as cellular rearrangements and tissue fluidity, to achieve a more realistic representation of morphology in dynamic tissues. Further, more systematic investigations of different extrusion mechanisms, such as basal protrusions (lamellipodia) [52], present another interesting avenue for future work.

Our analytical model does not fully describe the morphology of partially extruded cells because we neglected the relaxation of tilt in the neighboring cells, i.e., a nonvanishing angle between the tissue-midplane normal and lateral interfaces. To better capture such intermediate morphologies of cell extrusion, as observed numerically in the BVM, and thus also obtain better predictions of stability, it is necessary to find a continuum description for these states. One possibility would be to introduce a continuous tilt field, to allow for deviations in the lateral configurations. Such a field has been introduced before in the context of lipid membranes [55,56]. Tilt has been linked to scutoid cell shapes, which should be investigated in this context as well [57,58]. Together with the framework of a discrete cell in a mean-field tissue, which we presented here, it can serve as a starting point for further investigations of the coupling of single-cell and tissue-scale mechanics.

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DATA AVAILABILITY

Both the computer code [33] and the data [59] that support the findings of this article are openly available.

APPENDIX A: DESCRIPTION OF ORGANOIDCHASTE

The vertex model is implemented in *OrganoidChaste* [33], which is a package for Chaste software [34]. Cells are polyhedra with apical, basal, and lateral faces, which are described by nodes in 3D; cf. Fig. 1(a). As nodes of faces are not necessarily coplanar, we interpolate the interfaces by considering an average node for triangulation [19,30]. To minimize the energy, Eq. (1), the force on node i is calculated based on the energy gradient with respect to the positions of node $\mathbf{r}_i$,

$$\mathbf{F}_i = -\nabla_{\mathbf{r}_i} E. \quad (\text{A1})$$

Overdamped dynamics is assumed in the equation of motion of the nodes,

$$\frac{d}{dt} \mathbf{r}_i = \mathbf{F}_i, \quad (\text{A2})$$

where we assumed a mobility of 1. This equation of motion is solved using Euler stepping. Volume conservation is achieved by projecting the forces on the space of volume conserving forces via a Lagrange multiplier, corresponding to pressures, with an additional correction step. This implementation is identical in our VM code and in the subsequent BVM minimization in Surface Evolver [36] described below.

Cells can perform neighbor exchanges, so-called T1 transitions, where we first form a protorosette (four-way one-dimensional junction) and then, in the next step, create the new interface, following the approach suggested in Ref. [30]. We perform these randomly with a probability given by a corresponding transition rate on the entire epithelium. This approach has been found to fluidize the 3D VM, allowing for a viscous relaxation of stresses [30]. In the minimization of the spherical epithelia, a decreasing transition rate is taken, similarly to simulated annealing. The resulting organoid is taken as the starting point for a comparative minimization in the VM and the BVM with the same underlying topology.

In order to solve the bubbly vertex model the configuration that is obtained in *OrganoidChaste* is exported to Surface Evolver [36], where the same energy, Eq. (1), is

considered. However, the interpolation points of the triangulation are no longer averaged, allowing for interfacial curvature. Nevertheless, the model and the minimization through overdamped relaxation are identical to the VM if no additional effects are added; in particular, the energy is unchanged, Eq. (1). If an interfacial bending rigidity is considered, then the mean curvature in Surface Evolver is calculated at triangulation points of the interfaces, which do not belong to the edges of three intersecting interfaces. This process leads to small deviations at the corner triangles due to the neglect of a bending energy across the corresponding interface internal edge [cf. Fig. 4(c)], which decreases as the triangulation resolution is increased.

Simulations of the VM were run on the timescale of minutes on a regular computer. Minimizations in Surface Evolver for large systems were run on a high-performance computing cluster on AMD Epyc 7H12 processors, with a minimization running on one 2.6-GHz core for a time from a few hours to up to three days for very large systems; cf. Fig. 5(c).

APPENDIX B: EXPERIMENTAL METHODS

Mouse colon organoids were generated from a male C57BL/6J mouse as described in Ref. [38]. To maintain organoid cultures, mCol organoids were enzymatically and mechanically dissociated into single cells using TrypLE (Gibco). Cells were resuspended in Basement Membrane Extract Type R1 (BME-R1) (R&D Systems), embedded in 10- μL domes in wells of 12-well plates (Corning), resulting in 5 domes per well. Organoids were cultivated in WENR medium prepared by the addition of 1 $\times$ B-27 supplement (Gibco), 1.25 mM n-Acetyl Cysteine (Sigma-Aldrich), 100 mgmL^{-1} Primocin (InvivoGen), 50 ngmL^{-1} EGF (Peprotech), 10% (v/v) Noggin Conditioned Medium, 20% (v/v) R-spondin Conditioned Medium, 10 μM Y-27632 (TargetMol), and 0.5 nM Wnt Surrogate (Gibco) into AD3+ basal medium, which consisted of Advanced DMEM/F-12 (Gibco) supplemented with 10 mM HEPES (Sigma-Aldrich), 2 mM GlutaMAX (Gibco), and 10 mM Penicillin/Streptomycin (Sigma-Aldrich). Organoids were cultured at 37°C in a humidified atmosphere with 5% CO_2 (Binder). The culture medium was regularly changed every two to three days, and passaging was performed at a 1:6 ratio approximately every five to seven days.

The mCol organoids were dissociated into single cells as described above and seeded at a density of 1000 cells μL^{-1} BME-R1. The organoids were embedded in 10 μL BME-1 domes in 1 well of 8-well μ -slide (ibidi) and cultured at 37°C in a humidified atmosphere with 5% CO_2 . The culture medium was changed on the second day of culture. For Latrunculin-A perturbation, organoids were treated for 5 hours with 0.5 μM Latrunculin-A (Cayman Chemical). In the corresponding paired experiment, the wild-type organoids were left in WENR medium. Following treatment, organoids were fixed (4% formaldehyde) for

30 minutes at room temperature. 0.5 μM Latrunculin-A was added to the fixation solution for organoids treated with Latrunculin-A. For immunofluorescence staining, organoids were permeabilized in 0.5% Triton-X100 in PBS (Sigma-Aldrich) for 10 minutes at room temperature. Organoids were subsequently incubated in blocking buffer (4% BSA, 0.1% Triton X-100, and 0.1% Tween-20 in PBS) for 1 hour at room temperature. The organoids were then incubated overnight at 4°C with primary antibody (Mouse anti-Beta-Catenin, BD Biosciences, 1:500). Organoids were washed with 0.1% Tween-20 in PBS 3 times before secondary antibody incubation for 1 hour at room temperature (Goat-anti-Mouse-AF488, Life Technologies, 1:500). Costaining for F-actin (Phalloidin-TRITC, Sigma, 50ngmL⁻¹) and DNA (Hoechst, Life Technologies, 1 μgmL^{-1}) was performed simultaneously. After three additional washes with 0.1% Tween-20 in PBS, the organoids were immersed in clearing solution (2.5 M fructose in 60% glycerol).

Images were acquired on a Nikon AX laser scanning confocal microscope using a 25 $\times$ silicone immersion objective, with an XY pixel size of 0.25 μm and a z-step size of 0.5 μm . Images were deconvolved using the automatic 3D deconvolution feature in NIS Elements software Version 5.42.03.

APPENDIX C: IMAGE SEGMENTATION

For image segmentation, the images were preprocessed by interleaving the image stack with averaged images to obtain cubic voxels. Further, the pixel intensity was linearly corrected as a function of imaging depth for some organoids, to correct for intensity loss. Finally, in some organoids, gamma correction was performed to increase the image contrast. Using CELLPOSE2.0 [60,61] and ilastik [62], with the models trained specifically for organoid imaging data with mouse small intestinal organoid data in Ref. [39], we segmented the spherical mouse colon organoids.

Using the force inference pipeline we have developed [39], we then performed three-dimensional reconstruction of the interfaces to obtain normal vectors at the tricellular junctions and subsequent force inference. For this process, we assumed constant surface tensions across cell-cell, cell-lumen, or cell-outside interfaces, which implies the force balance equation at a junction [cf. Fig. 2(c)],

$$\begin{aligned}\Gamma_1 + \Gamma_2 \cos(\theta_{12}) + \Gamma_3 \cos(\theta_{13}) &= 0, \\ \Gamma_2 \sin(\theta_{12}) - \Gamma_3 \sin(\theta_{13}) &= 0,\end{aligned}\quad (\text{C1})$$

where the three interfaces have tensions $\Gamma_{1,2,3}$, respectively, and the angle between interface i and j is denoted by θ_{ij} . Including all interfaces and all junctions yields a large overdetermined system of equations, where the angles can be used to infer the underlying (relative) surface tensions Γ [39], with the average tension over all interfaces set to 1. As

in Ref. [39], interfaces that could be associated with segmentation errors were excluded from the analysis.

Both in the BVM and in the experimental data, we determined the opening angle of a lateral interface by calculating the angle between the apicobasal axis [the vector pointing from the apical (luminal) interface center of mass to the basal (outer) one] and the inward pointing lateral interface normal φ via $\delta/2 = \pi/2 - \varphi$. For each cell, these lateral opening angles were averaged. For this calculation, interfaces were represented as point clouds. In the BVM, we considered the triangulation points. For the experimental data, we obtained the points by considering voxels that have two cell indices in their 19-point stencil (next-to-nearest neighbor voxels) and taking their positions in voxel space. We determined the lateral normals as the direction with smallest variance via a principal component analysis [39].

To determine the number of neighbors for each cell, we considered the basal (outer) side, as we found that segmentation errors usually occurred on the inner apical side. In the segmentation, we sometimes found inclusions of small cells due to segmentation issues [39]. We thus neglected cells with a volume below 1000 voxels (15.625 μm^3) and then determined tri-interfacial junctions, again via the 19-point stencil. For each basal interface, the number of neighboring junctions, i.e., polygonal edges, was determined as the number of neighbors. Cells in which segmentation errors were identified (as in the force inference step) were then neglected in the statistics. Their basal sides, however, have contributed to the neighbor determination, thus not leading to an error from this correction.

APPENDIX D: ESTIMATING LUMINAL PRESSURE

To estimate luminal pressures, we considered the vertex model with flat interfaces and its coarse-grained nonlinear stretching energy density [19]

$$e_{\text{stretch}} = \frac{\Gamma_a + \Gamma_b}{\Gamma_l} \Gamma_0 \left(\lambda^2 + 2 \frac{1}{\lambda} - 3 \right), \quad (\text{D1})$$

with average tension $\Gamma_0 = 5 \text{ nN}\mu\text{m}^{-1}$ and principal stretches in the spherical shell $\lambda = \lambda_\phi = \lambda_\theta$, which describe an elongation $\ell' = (1 + \lambda)\ell$ for an in-plane length ℓ that is stretched to ℓ' . One can show that bending energy contributions constitute higher-order corrections, which can be neglected.

To obtain the pressure, we considered

$$Pd[V(R)] = d[e_{\text{stretch}}(R)A_0], \quad (\text{D2})$$

with pressure P , spherical volume V , and undeformed spherical area A_0 ; the latter two both depend on the pressured sphere radius R . For a stretched sphere, the principal stretch was $\lambda = R/R_0$ with undeformed sphere

radius R_0 . We can insert the volume and area of a sphere to obtain

$$P4\pi R^2 dR = 4\pi \frac{\Gamma_a + \Gamma_b}{\Gamma_l} \Gamma_0 \left(2R - 2\frac{R_0^3}{R^2} \right) dR, \quad (\text{D3})$$

which implies, for the pressure,

$$P = \frac{\Gamma_a + \Gamma_b}{\Gamma_l} \Gamma_0 \left(\frac{2}{R} - 2\frac{R_0^3}{R^4} \right). \quad (\text{D4})$$

Note that the first term is the Young-Laplace law, which is modified by the second term stemming from the initial shell radius. For the shell in Fig. 2(a) with $\Gamma_a \approx 1.02$, $\Gamma_b \approx 1.06$, $\Gamma_l \approx 0.97$, and $R \approx 22 \mu\text{m}$, we have $R_0 = (\psi_{\text{exp}}/\psi_{\text{theo}})^{1/3} R \approx 18 \mu\text{m}$, which implies $P \approx 440 \text{Pa}$.

APPENDIX E: EFFECT OF LINE TENSION

The existence of line tensions has been demonstrated in small intestine organoids [39] and could, in principle, stabilize cell shapes basally. To consider such a process in our model, we added a line tension at all outer basal interface boundaries, which effectively constitutes a contractile contribution along the basal interface edges. Figure 10 displays the opening angles of BVM simulations with basal (outer) line tensions for all cells, where we see cell-specific variability in switching opening angles through an unbulged intermediate step. The cell-specific switching again suggests that cell-to-cell variability in tensions is necessary.

Besides a reduction of apicobasal interfacial tension, other possible mechanisms for extrusion are contractile actin structures pushing cells outward [3]. To introduce these mechanisms, we consider a line tension on the (luminal) apical side. For a tension τ , we thus have an energy at the defect

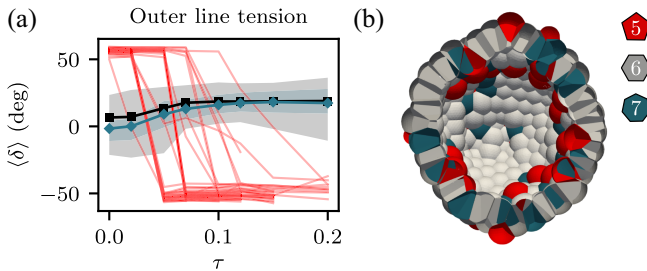


FIG. 10. Effect of basal (outer) line tensions in all cells on the opening angles of defect cells. (a) Opening angles as a function of the dimensionless line tension τ . Individual red lines mark single pentagonal cells to show single-cell dependency of inward bulging. Symbols show means, with shaded regions indicating 5% and 95% quantiles for hexagons in black and heptagons in blue. (b) Spherical shell as in Fig. 1 for $\tau = 0.05$.

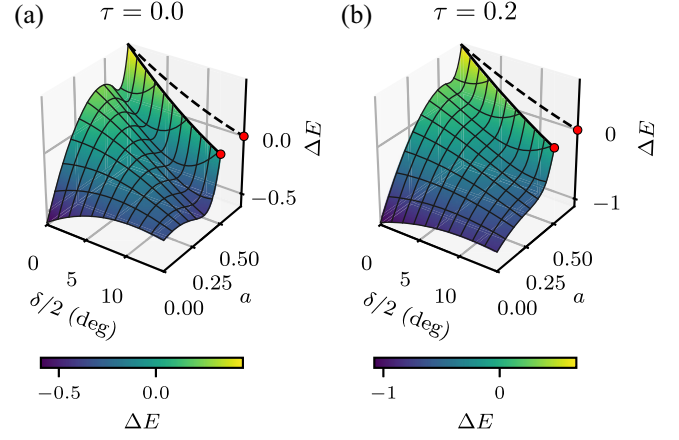


FIG. 11. Energy landscapes in configuration space for a pentagonal defect in a tissue with curvature screening and basal line tension τ . The cellular opening angle $\delta/2$ is bounded from above because of the maximum opening angle of the cap; see Fig. 9(a). Color-coding and lines are as in Figs. 8(a)–8(c). Parameters are $\Gamma = 0.6$, $R_{\text{bg}}/\bar{h} = 5$.

$$E_{\text{line}} = n\tau a, \quad (\text{E1})$$

which is an additional a -dependent term contributing to the total energy. It leads to a tilting of the energy landscape in configuration space because, in the case of curvature screening, no R_0 minimization is performed. Figure 11 displays such energy landscapes for pentagons ($n = 5$) for different line tensions τ . We see that the tilting can destabilize bulged states. However, since at $a = 0$ the line tension does not contribute anymore, an energy barrier in the case of a transition from a bulged state with $a = 0$ to the extruded case is not changed by the line tension. In light of partially extruded cells, such as in Fig. 8(g), this result implies that an additional tuning of interfacial tensions (either lateral or apicobasal) is necessary to achieve full extrusion.

APPENDIX F: DETAILS ON MEAN-FIELD TISSUE PARAMETRIZATION

The deformation energy for a tissue parametrized as in Fig. 6 comprises bending and stretching components, derived in the following. For the cap segment, we consider the radial deflection function f with rotational symmetry

$$f(r) = R_0 - \sqrt{R_0^2 - (\rho + r)^2}, \quad (\text{F1})$$

where ρ is the radial offset due to the bulged defect cell; cf. Fig. 6. The bending energy density reads

$$e_{\text{bending}} = \frac{\kappa}{2} H^2 + \kappa_G K, \quad (\text{F2})$$

with bending rigidity κ , saddle splay modulus κ_G , (total) mean curvature $H = c + c'$, and Gauss curvature $K = cc'$ for principal curvatures c, c' . For the deflection function $f(r)$, the curvatures read

$$H(r) = \frac{1}{r} \partial_r r \left(\frac{\partial_r f}{\sqrt{1 + (\partial_r f)^2}} \right) = \frac{2}{R_0} + \frac{\rho}{r R_0},$$

$$K(r) = \frac{\frac{1}{r} (\partial_r f) (\partial_r^2 f)}{(1 + (\partial_r f)^2)^2} = \frac{\rho + r}{r R_0^2}. \quad (\text{F3})$$

With the area element, $dS = \sqrt{1 + f'^2} r d\phi dr$, the corresponding bending energy from the mean curvature reads

$$\frac{E_{\text{bending}}^{\text{mean}}}{\kappa\pi} = -\frac{4}{R_0} \sqrt{R_0^2 - (\rho + r)^2} - \frac{\rho^2}{\sqrt{R_0^2 - \rho^2}} \times \text{artanh} \left(\frac{R_0^2 - \rho(\rho + r)}{(R_0^2 - \rho^2)^{1/2} (R_0^2 - (\rho + r)^2)^{1/2}} \right) \Bigg|_{r=r_0}^{r=R_B}. \quad (\text{F4})$$

Note that, for $\rho = r_0 = 0$ and small buckling radius R_B , i.e., small $|\nabla f|$ in the outer cone segment, the Taylor approximation in R_B scales quadratically, consistent with the approximation of the total energy with a flat disc as proposed by Seung and Nelson [23] for the flattened tip for the deformed shape around the defect. The first term in Eq. (F4) arises from the spherical contribution of the mean curvature bending energy; cf. corresponding first term in the mean curvature in Eq. (F3). The second term, which is proportional to the squared radial offset ρ^2 , serves as a correction for reducing the radius of the circular segment in cylindrical coordinates, as we decrease the radius of curvature in the azimuthal direction via this offset. This term stems from the quadratic contribution of ρ^2 from the squared mean curvature, Eq. (F3), in the bending energy.

The bending energy from Gaussian curvature only depends on the opening angle of the defect cell δ , as this modifies the solid angle covered by the cap segment, consistent with the Gauss-Bonnet theorem,

$$E_{\text{bending}}^{\text{Gauss}} = 2\pi\kappa_G [\cos(\delta/2) - \sin(\gamma/2)], \quad (\text{F5})$$

with tissue opening angle $\gamma/2$; cf. Fig. 6(a).

The conical part of the bending energy is unchanged from the respective result for large $|\nabla f|$ derived in Ref. [23],

$$E_{\text{cone}} = \frac{f_1}{(1 + f_1^2)^{1/2}} \pi\kappa \log(R/R_B), \quad (\text{F6})$$

with the total radius R , the buckling radius R_B , and the cone slope

$$f_1^{\text{cone}} = \left[\frac{1}{(1 - s/2\pi)^2} - 1 \right]^{1/2}. \quad (\text{F7})$$

This slope can be obtained by considering a cone where the base's circumference corresponds to the circumference of the corresponding flat disc segment with a wedge of angle s cut out. In the case of a pentagonal defect, this implies $s = 2\pi/6$.

The stretching energy density reads

$$e_{\text{stretch}} = \frac{1}{2} (2\mu \varepsilon_{ij}^2 + \lambda \varepsilon_{kk}^2), \quad (\text{F8})$$

with two-dimensional Lamé coefficients μ and λ . To compute the stretching energy of the tissue, we consider the in-plane deformation of a conically bent elastic sheet. Using polar coordinates (r, θ) , it is clear from symmetry that the azimuthal deformation function must be $u_\theta = q_s \theta r$ because we have a stretch of $q_s r$ if we denote with q_s the topological charge (i.e., the relative material fraction of necessary azimuthal stretching). For a pentagonal defect, this means $q_s = s/2\pi = 1/6$ because $1/6$ of the azimuthal circumference is removed.

Introducing the radial position function $w(r)$, the radial deformation is $u_r(r) = w(r) - r$, and the strains for a (moderately) bent plane [63] in polar coordinates are described by

$$\varepsilon_{rr} = w'(r) - 1 + \frac{(f'(r))^2}{2},$$

$$\varepsilon_{\theta\theta}(r) = \frac{w(r) - r}{r} + q_s,$$

$$\varepsilon_{r\theta} = 0, \quad (\text{F9})$$

where the assumption of moderate bending yields a non-linear coupling term of the deflection function in the strain.

The energy containing both stretching and bending contributions [cf. Eq. (9)] hence constitutes a functional for w . The corresponding Euler-Lagrange equation for w is

$$\frac{w}{r} - w' - r w'' + \frac{2\mu}{2\mu + \lambda} \left(q_s - \frac{(f')^2}{2} \right) - r f' f'' = 0. \quad (\text{F10})$$

Determining the Euler-Lagrange equation for f from the sum of bending and stretching energy and considering a cone $f' = \text{const}$ yields the identical result as obtained in Ref. [23], implying the bending energy $E = \pi\kappa \log(R/r_0)$ for some cutoff r_0 . For this result from Ref. [23], one uses $w(0) = 0$ and $\sigma_{rr}(R) = 0$ as boundary conditions.

We now want to take a different approach, however, and consider a ring of material with inner and outer radii r_0 and R_B , deflected with a known fixed slope $f' = f_1$. Introducing $\Delta s = q_s - f_1^2/2$, Eq. (F10) is solved by

$$w(r) = \frac{\mu}{2\mu + \lambda} \Delta s r \log(r) + \frac{k_1}{r} + k_2 r, \quad (\text{F11})$$

with coefficients k_1 and k_2 , which are determined from the boundary conditions.

Note that the strains in Eq. (F9) only consider the leading-order terms in f' . To also allow for large f' , we observe that the azimuthal stretch in a cone is reduced such that, for a circle at radius r , the deformed radius becomes $r/\sqrt{1+f_1^2}$. This results in the corrected radial strain $\varepsilon_{rr} = w'(r) - (1+f_1^2)^{-1/2}$. Since we only consider a linear f , i.e., we neglect f'' , only the term corresponding to Δs changes in Eq. (F10), implying that the solution is again Eq. (F11), now with

$$\Delta s = q_s - \left(1 - \frac{1}{\sqrt{1+f_1^2}}\right). \quad (\text{F12})$$

As boundary conditions, we assume vanishing radial stresses at the inner and outer radii, i.e., at $r = r_0$ and $r = R_B$,

$$\sigma_{rr}(r) = 2\mu\varepsilon_{rr}(r) + \lambda(\varepsilon_{rr}(r) + \varepsilon_{\theta\theta}(r)) = 0. \quad (\text{F13})$$

This boundary condition assumes that no radial (and thus total compressive) stresses act on the boundary of the cap segment and the conical part of the surrounding tissue. Comparing the cell height in VM simulations to the theoretical mean-field tissue height serves as an estimate for the compressive in-plane stress. The height difference from the mean-field height (cf. Appendix H) is shown in Fig. 12(a).

We see that our parametrization for the tissue shape seems reasonable in comparison to the VM shape: Around the defect cell, we have a ring of tissue that interpolates the opening angle of the cell and the tissue [corresponding to the light gray tissue in Fig. 6(a)], surrounded by a conically deformed background tissue [dark gray tissue in Fig. 6(a)]. The height difference decays toward the boundary between the cap and the cone segments with a vanishing height difference (and thus stress) in the outer cone segment. Note, however, that our assumption on the defect cell height does not perfectly match the simulations. An additional coupling of tissue stress at the inner ring of the surrounding tissue and the defect cell's slant height could be introduced. We neglect this additional degree of freedom here in order to obtain a tractable model, and we obtain good agreement to simulations nonetheless.

After determining the coefficients $k_{1,2}$, we integrate the stretching energy to find

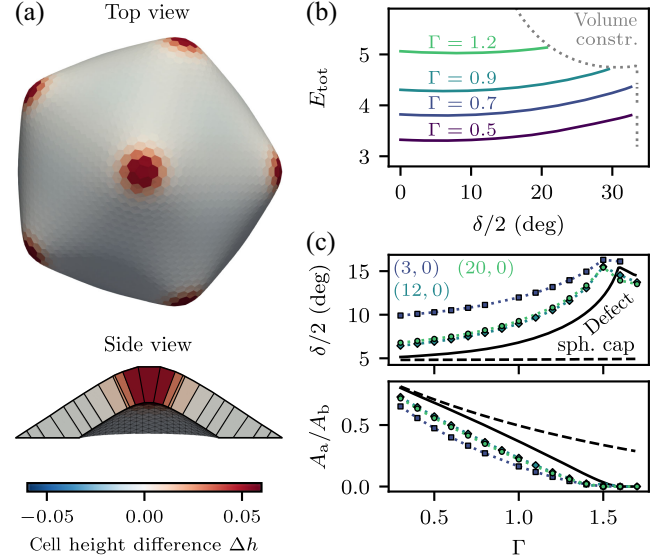


FIG. 12. Comparison of the continuum model with vertex model simulations of icosahedral shells. (a) VM shell with Caspar-Klug size (20, 0) and $\Gamma = 1.1$ with top-view and side-view cross sections. The color-coding shows the cell height difference from the mean-field VM height [cf. Eq. (H1)]. (b) Total energy of the proposed full defect model as a function of defect opening angle $\delta/2$ for different apical-basal surface tensions Γ . The dotted gray line indicates the angle boundary from the volume constraint. (c) Opening angles $\delta/2$ (top) and apical-to-basal face area ratios (bottom) for the different models. The defect model (solid line) considers a defect in a mean-field tissue with the angle $\delta/2$ corresponding to the minimal energy configuration; cf. panel (a). The spherical cap model (dashed line) corresponds to this energy under the assumption that $\delta/2$ matches the background curvature. The symbols indicate vertex model simulation results for different Caspar-Klug indices (T_1, T_2) .

$$\frac{E_{\text{stretch}}}{\pi Y(1+f_1^2)^{1/2}} = \frac{\Delta s^2}{8} (R_B^2 - r_0^2) - \frac{\Delta s^2}{2} \frac{R_B^2 r_0^2}{R_B^2 - r_0^2} \log^2 \left(\frac{R_B}{r_0} \right). \quad (\text{F14})$$

Coming back to the mean-field tissue, we now approximate the stretching energy in the cap segment [light gray in Fig. 6(a)] by replacing it with a cone segment such that the slope matches the opening angle of the defect cell, $f_1 = \tan(\delta/2)$. If we consider the fully extruded case with $r_0 = 0$ and $\delta = 0$, we have the well-known analogous version of a flat disclination for large $|\nabla f|$ [23]. The corresponding approximation for small slopes has been used successfully in the description of viral capsids [20]. For the conical case, where the full Gaussian curvature is relaxed in the defect cell and $f_1 = f_1^{\text{cone}}$, we have $\delta/2 = \pi/6$ and $\Delta s = 0$, as expected. Our approximation thus serves as an interpolation between the two extreme limits of the stretching energy of a flat disclination (as used for viral

capsids) and the solution for a cone, depending on the opening angle of the defect cell.

For the outer conical segment of the mean-field tissue [dark gray in Fig. 6(a)], we assume that it attains the conical shape, and thus the stretching energy vanishes.

Finally, we can write down the full energy of the surrounding tissue as a sum of the contributions in Eqs. (F4)–(F6) and (F14), i.e.,

$$E_{\text{tissue}} = E_{\text{bending}}^{\text{mean}} + E_{\text{bending}}^{\text{Gauss}} + E_{\text{cone}} + E_{\text{stretch}}. \quad (\text{F15})$$

APPENDIX G: VERIFICATION OF DEFECT CELL TREATMENT IN ICOSAHEDRAL SHELLS

Because of the analytical tractability, icosahedral shells constitute a good test case. In Ref. [19], we have shown that their shape follows the described theory of moderately bent plates. Comparing to simulations, we now show that the above parametrization and continuum theory are able to describe the shape of the pentagonal cell and of the surrounding tissue found in the VM.

The corresponding continuum model energy for this vertex model case consists of the core energy for a defect cell without bulging apical-basal membranes, i.e., with $h_s^a = 0 = h_s^b$, Eq. (8); the bending energy of both the cap segment, Eqs. (F4) and (F5), and the outer cone region, Eq. (F6); and the stretching energy in the cap, Eq. (F14). For the resulting total energy E_{tot} , we assume a large outer radius of the total disc R , such that area changes from deformations in the cone are negligible; thus, we drop the term proportional to $\log(R)$ in Eq. (F6), as we look at energy differences with respect to a reference state.

In the derivation of the mean-field tissue (cf. Appendix H) lateral interface energies in the energy density are included by counting half of them for the respective cell area (to avoid double counting). Here, we introduce a boundary energy term, which subtracts this lateral energy from the mean-field tissue and adds half the pentagonal lattice energy of the defect. This lateral energy is present in full and should be considered completely due to the explicit description of the defect cell. For the lateral contact between the defect and the tissue, we thus subtract the lateral energy contribution that is implicitly given in the mean-field tissue boundary (with a circular cutout) and replace it with the actual lateral energy contribution of the cell (with a pentagonal shape). This energy reads

$$E_{\text{bdry}} = \frac{5\tilde{h}}{4} \left(2a + \frac{\tilde{h}}{\eta_{\text{geom}}} \sin(\delta/2) \right) - \pi r_0 \tilde{h}, \quad (\text{G1})$$

and it is added to the total energy, i.e.,

$$E_{\text{tot}} = E_{\text{core}} + E_{\text{bending}}^{\text{mean}} + E_{\text{bending}}^{\text{Gauss}} + E_{\text{cone}} + E_{\text{stretch}} + E_{\text{bdry}}. \quad (\text{G2})$$

For the elastic parameters and the bending rigidities of the mean-field tissues, we use the mean-field values derived for the VM in Ref. [19], i.e., Eqs. (H2) and (H4) in Appendix H.

We find that for every $\delta/2$ and a , the energy has a unique minimum as a function of the buckling radius R_B or, equivalently, as a function of the radial cap curvature radius R_0 . We thus minimize the energy with respect to R_B to obtain the total energy as a function of the opening angle $\delta/2$ in the vertex model. Figure 12(b) depicts the energies for different surface tensions Γ , where we observe a clear minimum. Note that, for large Γ , the volume constraint yields an upper bound for the possible opening angles. Our approach, the defect model, allows the defect to have an opening angle independently from the background tissue curvature. For comparison, we also consider the case of the opening angle following the background curvature; i.e., $\delta/2$ is chosen such that it matches R_0 with $\rho = 0$, which we henceforth denote as the spherical cap model. The offset of the energies is a consequence of neglecting the $\log(R)$ term, and it results from the different cell number densities for different Γ .

In Fig. 12(c), we show the opening angle and the ratio of basal and apical interfacial areas in the minimal energy configuration. We see that the defect model predicts a strong increase of $\delta/2$ until $\Gamma \approx 1.4$ when the basal area completely collapses and the angle is determined by the volume constraint. The ratio of basal and apical areas also shows this collapse. The symbols indicate vertex model simulations of spheres with different Caspar-Klug indices (T_1, T_2). Our defect model agrees well with the numerical vertex model data for large enough spheres and captures the shape of the defect much better than the model without the defect's ability to adopt a different opening radius as the surrounding tissue. Note that we do not have any fitting parameters but rather derive the full model from the vertex model with only one free parameter, namely, Γ . We thus conclude that the approach of considering a single defect cell in a mean-field tissue captures the resulting shape better than an approach purely based on continuum theory, as exemplified by the comparison to the VM here.

APPENDIX H: CONTINUUM PARAMETERS FOR THE MEAN-FIELD TISSUE

For the continuum description of the VM, we consider a flat prototypical hexagonal cell without interfacial curvature, described by the VM energy, Eq. (1). Because of volume conservation, $V = 1$, we can relate the hexagonal edge length and the height. Minimizing the energy with respect to height yields the equilibrium height [19]

$$h = 2^{1/3} 3^{-1/6} (\Gamma_a + \Gamma_b)^{2/3}. \quad (\text{H1})$$

To derive the elastic stretching parameters, in Ref. [19], a single-cell deformation is considered. For the same energy and hexagonal cells, the Lamé coefficients read

$$2\mu = \lambda = (\Gamma_a + \Gamma_b), \quad (\text{H2})$$

which implies, for the 2D Young's modulus and Poisson ratio,

$$Y = \frac{4\mu(\mu + \lambda)}{2\mu + \lambda} = \frac{3}{2}(\Gamma_a + \Gamma_b),$$

$$\nu = \frac{\lambda}{2\mu + \lambda} = \frac{1}{2}, \quad (\text{H3})$$

respectively. To derive the bending rigidity and saddle splay modulus, a curved cell is considered in Ref. [19], and the energy change as a function of the cellular curvature for small curvatures is determined to find

$$\kappa = \frac{9}{8} \frac{1.26}{2^{1/3} 3^{4/3}} (\Gamma_a + \Gamma_b)^{1/3},$$

$$\kappa_G = \left[\frac{(\Gamma_a + \Gamma_b)^2}{2} + \frac{3}{2} - \frac{9}{4} 1.26 \right] \frac{(\Gamma_a + \Gamma_b)^{1/3}}{2^{1/3} 3^{4/3}}. \quad (\text{H4})$$

In Ref. [19], it has been shown that the VM does not perfectly follow linear elasticity and that the stretching moduli are functions of the strains. Thus, a nonlinearity correction for the icosahedral instability is necessary when considering the continuum limit of the VM, i.e., $k_{\text{ico}} = 1/2$.

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